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Simulating the chemical kinetics of CO₂-methane exchange in hydrate

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Abstract

Carbon dioxide exchange with methane in the clathrate structure has been shown beneficial in laboratory experiments and has been suggested as a field-scale technique for production of natural gas from gas-hydrate bearing sediments. Furthermore, the method is environmentally attractive due to the formation of CO₂-hydrate in the sediments, leading to the geosequestration of carbon dioxide. However, the knowledge is still limited on the impact of small-scale heterogeneities on hydrate dissociation kinetics. In the present study, we developed a model for simulating laboratory experiments of carbon dioxide injection into a pressure vessel containing a mixture of gas hydrate and quartz sand. Four experiments at different temperature and pressure conditions were modeled. The model assumes that the contents are ideally mixed and aims to estimate the effective dissociation rate of gas hydrate by matching the model results with the experimental observations. Simulation results indicate that with a marginal offset the model was able to simulate different hydrate dissociation experiments, in particular, those that are performed at high pressures and low temperatures. At low pressures and high temperatures large discrepancies were noticed between the model results and the experimental observations. The mismatches were attributed to the development of extremely heterogeneous flow patterns at pore-scale, where field-scale models usually assume the characteristics to be uniform. Through this modeling study we estimated the irreversible dissociation rate of methane- and CO₂-hydrate as 0.02 and 0.03 mol.m⁻³s⁻¹, respectively.

Keywords: *CO₂ injection; CO₂-methane exchange; Gas-hydrate recovery; Small-scale heterogeneities; Kinetic modeling*

23 1 Introduction

24 Gas-hydrates are solid clathrate compounds that are thermodynamically stable at low temper-
25 atures and high pressures. Such conditions naturally exist below permafrost and in deep ocean
26 sediments in which immense amount of methane is estimated to be stored as gas-hydrate de-
27 posits (Archer et al., 2009; Burwicz et al., 2011). The global amount of gas-hydrate deposits
28 have been reported between 10^{15} and 10^{18} standard cubic meters (Piñero et al., 2013; Wallmann
29 et al., 2012), or about 15 Tera tonnes of oil equivalent (Makogon, 2010) which is adequate for
30 maintaining the supply of energy for centuries. Although the range of estimates is wide, it is
31 agreed that the available amount of gas-hydrate deposits is huge and thus worth of the atten-
32 tion as an alternative source of energy. Development of strategies for extraction of methane
33 from gas-hydrate reservoirs has recently become an economically attractive option given the
34 environmental desirability of natural gas as a fuel in comparison to other fossil fuels.

35 Methods of producing natural gas from gas-hydrates are mainly based on disturbing the ther-
36 modynamic stability of gas-hydrate in the reservoir leading to dissociation of the gas-hydrate
37 and release of the methane. The methods include (i) thermal stimulation by increasing the tem-
38 perature in the reservoir (e.g., Fitzgerald and Castaldi, 2013), (ii) depressurization (e.g., Ahmadi
39 et al., 2007), (iii) hydrate conversion by substituting gas molecules inside the gas-hydrate crystals
40 with another similar gas (e.g., Kvamme et al., 2007, 2016; Ohgaki et al., 1996), and (iv) injection
41 of thermodynamic inhibitors (e.g., amino acids, salts, alcohols or non-ionic surfactants) (Erfani
42 et al., 2017; Masoudi and Tohidi, 2005) for altering phase equilibrium conditions. Amongst all
43 these methods, the conversion of methane-hydrate to CO_2 -hydrate by injection of CO_2 has par-
44 ticularly attracted attentions since carbon dioxide is shown to be able to displace methane in the
45 hydrate lattice provided that both gases form a similar hydrate structure (type SI) (Kvamme
46 et al., 2016; Ohgaki et al., 1996; Voronov et al., 2014). The replacement of guest molecules can
47 happen either directly without dissociation of the hydrate structure or indirectly through con-
48 secutive dissociation of methane-hydrate and formation of CO_2 -hydrate. Goel (2006) discussed
49 that the introduction of carbon dioxide to the reservoir and its conversion to hydrate is alone suf-
50 ficient to thermodynamically maintain the dissociation of methane-hydrate. The CO_2 -methane
51 exchange, regardless of its exchange mechanism, is particularly interesting for its capacity to
52 sequester carbon dioxide in favor of reducing greenhouse gas emissions (see e.g., Dashti et al.,
53 2015; Kvamme et al., 2007). The method also has a couple of other side benefits such as main-

54 taining the mechanical stability of the reservoir preventing sea-floor landslides in field operations
55 (Sultan et al., 2004), and the potential for thermal stimulation through the injection of super-
56 critical carbon dioxide (Deusner et al., 2012; Ebinuma, 1993). The feasibility of CO₂-methane
57 exchange as a technology to produce natural gas from gas-hydrate zones has already been pro-
58 posed and investigated (e.g., Yonkofski et al., 2016). Many other studies, e.g., Kvamme et al.
59 (2016); Deusner et al. (2012); Ota et al. (2005), analyzed the outcome of CO₂-methane exchange
60 at laboratory scale using apparatuses in which carbon dioxide (either gas or liquid) is injected
61 into a vessel containing methane-hydrate. A substantial number of studies have used numerical
62 models to evaluate the conventional methods of production from gas-hydrate reservoirs (e.g.,
63 Moridis and Reagan, 2011a,b; Vafaei et al., 2014). However, numerical studies on CO₂-methane
64 exchange are few and are mostly limited to the field-scale. For example, White et al. (2011)
65 modeled the injection of carbon dioxide into a depressurized gas-hydrate reservoir and stated
66 that the low injection pressures of carbon dioxide can enhance the methane recovery from class
67 1 hydrate.

68 Although significant research efforts have been dedicated to the development of efficient ex-
69 perimental procedures and reliable models[they may ask for references], the complex reaction
70 kinetics of CO₂-methane exchange at the scales of pore to core has not yet been addressed in
71 detail or experimentally constrained under the controlled conditions. Most of current modeling
72 approaches [e.g. ???] simplify the reaction kinetics (usually employ a simple first-order kinet-
73 ics) and neglect the small-scale heterogeneities at the scale of their computational grid (where
74 the transport properties are averaged and considered constant).

75 In contrast to the existing modeling studies that mostly concentrated on complexity of
76 fluid dynamics at large scales (and simplified the reaction kinetics due to uprising numerical
77 instabilities), the present model focuses on complexity of the reaction kinetics and simplifies
78 the fluid flow mechanisms. To this end, the approach provides a measure to gauge the lone
79 importance of kinetics at small scales where heterogeneities are typically ignored. The overall
80 aim is thus to use the numerical simulations to unravel the extent of influence that typical
81 assumptions of simplifying reaction kinetics and ignoring pore-scale heterogeneities have on the
82 accuracy of estimations at small scales, and to illustrate the contributions of error to field-scale
83 modeling calculations. Are you sure about the word uprising above? Furthermore, the
84 present study evaluates the reported rate values of hydrate dissociation and formation in the

85 literature and approximates/testifies the effective rate parameter values for the experimental
86 results of [Deusner et al. \(2012\)](#). For this purpose, a rigorous optimization technique ([Babaei
87 and Pan, 2016](#)) was applied to fit the model to the experimental results.

88 The paper is structured as follows: first we describe the model structure and its underlying
89 assumptions. Then, the governing equations of hydrate dissociation/formation kinetics, mass
90 and energy balance are introduced. Next, we describe the optimization formulation to calibrate
91 the system kinetics using existing experimental data from [Deusner et al. \(2012\)](#). Finally, results
92 are presented and discussed.

93 2 Experimental Setup

94 [Deusner et al. \(2012\)](#) examined methane production from hydrates by injection of supercritical
95 carbon dioxide into a pressure vessel containing a water-saturated mixture of methane-hydrate
96 and quartz sand. The experiments were performed at four different pressure-temperature con-
97 ditions that are typical for naturally formed gas-hydrate reservoirs (Fig. 1).

98 The sediment samples were prepared at -20 °C from a homogeneous mixture of quartz sand
99 (grain size of 0.1-0.6 mm) and fine ice particles (grain size fraction of 0.3-1 mm) produced from
100 deionized water. Experiments were carried out in a custom-made high pressure stainless steel
101 apparatus. Supercritical CO₂ was injected with a piston pump from an inlet at the bottom of
102 the sample vessel and was heated to 95 °C inside temperature controlled conditioning chamber
103 prior to the injection. Pressure, salinity and temperature were continuously monitored and
104 recorded at the inlet and outlet. To achieve a constant rate of injection, pressure was adjusted
105 with a back-pressure regulator valve in line with a fine-regulating valve for the compensation of
106 pressure spikes. At the beginning of every CO₂ injection interval, the sediment-hydrate sample
107 was continuously percolated with saltwater at a flow rate of 1.0 ml.min⁻¹. The water pre-wash
108 was performed to ensure that the sample body was permeable and homogeneously pressurized.
109 CO₂ was injected stepwise following a sequential injection strategy and completed after a four
110 to six injection rounds with CO₂ supply rates of 2.5 to 5 ml.min⁻¹. The waiting time between
111 the injection intervals are referred to as equilibration intervals during which no effluent fluid was
112 produced and the system was left to reach thermodynamic equilibrium. During the equilibration
113 intervals, the system pressure was maintained by the injection of a small amount of CO₂ in order
114 to compensate the volume changes due to CO₂ cooling and phase changes. The CO₂ injection

115 intervals and the waiting time between them were different for each experiment.

116 Experiments were performed at three ambient temperatures (2, 8 and 10 °C). The temper-
117 ature was regulated at the exterior surface of the vessel with a thermostat system and kept
118 constant through the entire experiment. At the start of experiment, the vessel included only
119 three components: methane, water and quartz sand. Methane and water initially existed as
120 methane-hydrate. The quartz sand was assumed nonreactive and regarded as an inert solid
121 phase. During injection intervals, the introduction of hot CO₂ altered the system thermody-
122 namics and new additional components such as liquid CO₂ and gaseous methane were identified
123 (Fig. 2). CO₂-hydrate formation was also viable depending on system p/T conditions during
124 equilibration. It was impossible to exactly determine the final composition of gas-hydrate at the
125 end of the experiments. There is, however, a high possibility that a mixed CO₂-CH₄-hydrate
126 was formed in the vessel. Nevertheless, the mixed composition of gas-hydrates could not influ-
127 ence the mass balance calculations which were done based on component inventories and by the
128 volume balancing of inputs and outputs. See [Deusner et al. \(2012\)](#) for further details about the
129 experiments and the assembly of apparatus.

130 3 Materials and Methods

131 The model describes the experimental pressure vessel as an isobaric perfectly mixed reactor. In
132 this modeling approach, the system was considered homogeneous and the chemical components
133 inside it were assumed ideally mixed.

134 In the model, superheated liquid CO₂ entered from the inlet during the injection periods
135 and dissociates the methane-hydrate in place. Then the system is left to reach the equilibrium
136 and this cycle repeats for several times according to the experimental procedure. Given that
137 the vessel pressure was kept constant during the entire experiment, the mobile substances (e.g.,
138 water, CO₂ and methane) were allowed to discharge from the outlet during the injection intervals
139 only. The outlet composition was assumed identical to the composition of the substances inside
140 the reactor, which itself is a function of residence time and the reaction kinetics. Depending on
141 the p/T conditions in the vessel, CO₂- or methane-hydrate could form during the equilibrium
142 intervals. The terms CO₂-hydrate and methane-hydrate in this modeling study represent the
143 components CO₂ and methane incorporated in the gas-hydrate phase. The thermodynamics of
144 mixed gas-hydrates is not explicitly considered in the model.

145 3.1 Governing equations

146 3.1.1 Mass balance

147 System mass balance follows the equation (COMSOL 4.3, 2013),

$$\frac{d(c_j V_r)}{dt} = v_j c_{i,j} - v_j c_j + r_j V_r \quad (1)$$

148 where V_r is the volume of reactor, c_j is the concentration of substance j (Water, CO₂, methane,
149 etc.) in the system, $c_{i,j}$ is the concentration of substance j at the inlet, v_j is the rate of influent
150 stream to the system (equal to effluent) and r_j is the increase/decay rate of substance j according
151 to the reactions.

152 3.1.2 Energy balance

153 The solution of energy balance gives (COMSOL 4.3, 2013):

$$V_r \sum_j c_j C_{p,j} \frac{dT}{dt} = Q_r + Q_w + \sum_j v_j c_{i,j} (h_{i,j} - h_j) \quad (2)$$

154 where h_j is the enthalpy of substance j , $C_{p,j}$ the heat capacity of substance j , Q_w the amount
155 of energy lost or gained through the reactor walls, and Q_r the energy consumed or released by
156 reactions,

$$Q_r = V_r \sum_k H_k r_k \quad (3)$$

157 with H_k as the enthalpy of reaction k . Q_w is calculated analytically for the cylindrical shape of
158 reactor:

$$Q_w = 2\pi L \lambda_m (T_s - T) \quad (4)$$

159 where T is the system temperature, L the length of vessel, and T_s is the temperature at the
160 inner surface of reactor wall calculated by

$$T_s = \frac{\lambda_m T + \lambda_w T_w}{\lambda_m + \lambda_w} \quad (5)$$

161 with λ_m and λ_w calculated as,

$$\lambda_m = \frac{\kappa_m}{\ln(r_s/r_{in})}, \lambda_w = \frac{\kappa_w}{\ln(r_o/r_s)} \quad (6)$$

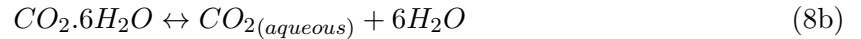
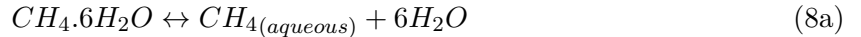
162 where r_s is the reactor inner radius, r_{in} is the radius of inlet, r_o is the reactor outer radius, κ_w is
 163 the thermal conductivity of the wall material and κ_m is the overall thermal conductivity of the
 164 system calculated by $\kappa_m = \sum s_j \kappa_j$ where κ_j is the thermal conductivity and s_j is the saturation
 165 of substances inside the vessel calculated by $s_j = c_j \phi M_j / \rho_j$. ϕ is the porosity of vessel, M_j is the
 166 molecular weight and ρ_j is the density of substance j . The enthalpy of substances at different
 167 system temperatures are calculated as,

$$h_j(T) = \int_0^T C_{p,j} dT + h_j(0) \quad (7)$$

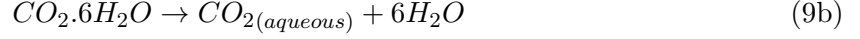
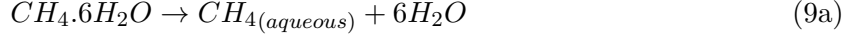
168 where $h_j(0)$ is the enthalpy of substance j at a reference temperature and pressure. $h_j(0)$ values
 169 at 293K and 13MPa for methane, CO₂ and water were calculated 12.25, 10.5 and 1.72 kJmol⁻¹
 170 respectively (NIST Chemistry WebBook, [Linstrom and Mallard, 2013](#)). $C_{p,j}$ was assumed con-
 171 stant for the p/T conditions of experiments.

172 3.1.3 Reactions

173 The solution of mass and energy balance considers the following reactions inside the reactor.
 174 Depending on the p/T conditions, hydrate dissolution and formation occur inside the hydrate
 175 stability region and hydrate dissociation occurs outside the hydrate stability region (see Fig. 1).
 176 The following pair of reversible reactions were considered under the stability conditions:



177 Reactions [8a-8b](#) account for hydrate dissolution while thermodynamically stable, but under-
 178 saturated with respect to the gas in the solution (water). Hydrate precipitation (formation)
 179 occurs at over-saturated conditions. A set of irreversible reactions were considered for p/T
 180 conditions at which hydrates are thermodynamically unstable,



181 Reactions 9a and 9b account for disintegration of hydrate when it is not stable. Since during the
 182 experiments the pressure of the system was kept constant, the stability of hydrates in model was
 183 determined only by the system temperature. The hydrate instability occurred when the system
 184 temperature exceeded the hydrate stability temperature T_c . For the experiments at 13MPa,
 185 the stability temperatures for CH₄-/CO₂-hydrate were measured from stability curves (Fig. 1)
 186 at 13.7/9.5 °C, respectively. These values were lower for the experiment at 8MPa and were
 187 determined to be 8.95 and 8.9 °C for CH₄- and CO₂- hydrate, respectively.

188 According to [Haeckel et al. \(2004\)](#) the rate of hydrate dissolution/formation r_j was calculated
 189 by,

$$r_j = \begin{cases} k_{rev,j}(c_{cte,j} - c_{Hydrate_j}) & \text{if } T < T_c, \\ k_{irr,j} & \text{if } T > T_c. \end{cases} \quad (10a)$$

$$r_j = \begin{cases} k_{rev,j}(c_{cte,j} - c_{Hydrate_j}) & \text{if } T < T_c, \\ k_{irr,j} & \text{if } T > T_c. \end{cases} \quad (10b)$$

190 Under hydrate stability condition, $k_{rev,j}$ is defined based on the Arrhenius formula,

$$k_{rev,j} = A_j e^{-\frac{\Delta E_j}{R_c T}} \quad (11)$$

191 where T is the system temperature, R_c is the universal gas constant, and for hydrate j , A_j
 192 denotes the frequency factor and ΔE_j the activation energy (Table 1). A_j is typically expressed
 193 in [mol.m⁻².s⁻¹.Pa⁻¹] and to convert its unit to [s⁻¹] the following equation is used ([Kim et al.,](#)
 194 [1987](#)):

$$A_j[s^{-1}] = \frac{6P_{sys}}{\Psi \rho_j d_j} A_j[mol m^{-2} s^{-1} Pa^{-1}] \quad (12)$$

195 where d_j is the average diameter of hydrate particles, P_{sys} is the system pressure, ρ_j is the hydrate
 196 molar density, and Ψ is the particles geometry term ($\Psi = 1$ for spherical particles). According
 197 to [Haeckel et al. \(2004\)](#), Eq. (10) assumes hydrate dissolution or formation to be proportional to
 198 the saturation of methane in pore water with respect to its equilibrium concentration (c_{cte}). At

199 hydrate instability conditions (at system temperatures above the stability temperature) $k_{irr,j}$
 200 was assumed constant and treated as an adjustable parameter.

201 The exchange rate of methane and CO₂ from pure phase to the water phase and vice-versa
 202 is defined by the following reversible reactions,



203 where the exchange rates follow the same mechanism as of Eq. (10) without the temperature
 204 dependencies, and similar to [Noyes et al. \(1996\)](#),

$$r_j = k_{s,j}(c_{cte,j} - c_j).$$

205 The exchange rate constants $k_{s,j}$ are estimated by the fitting procedure. The values of c_{cte}
 206 for aqueous CO₂ and methane, and both CO₂- and CH₄-hydrates at experimental temperatures,
 207 pressures and salinity are calculated according to Henry's law and listed in [Table 1](#).

208 3.2 Optimization technique

209 Interior-reflective Newton methods ([Coleman and Li, 1996](#); [Gharasoo et al., 2017](#)) which are
 210 often employed in solving optimization problems have difficulties in minimizing this model due
 211 to high nonlinearity and discontinuity of the objective function. We thus used a hybrid response
 212 surface surrogate-based method which also reduces the computation costs of the optimization
 213 process. The details of the algorithm is presented in [Babaei and Pan \(2016\)](#) where the authors
 214 showed that the surrogate model that most consistently and robustly results in a computationally
 215 efficient optimization operation is the Radial Basis Function (RBF).

216 We first define normalized root-mean-square derivations (NRMSD) for inventory CH₄ and
 217 CO₂ as functions of four unknown parameters $k_{irr,MGH}$, $k_{irr,CGH}$, $k_{s,CH4}$, and $k_{s,CO2}$:

$$\text{NRMSD}_{CO_2} = \frac{\sqrt{\mathbf{E}((CO_2^{inv.}) - (CO_2^{inv.exp}))^2)}}{\max(CO_2^{inv.exp}) - \min(CO_2^{inv.exp})} \quad (14a)$$

$$\text{NRMSD}_{CH_4} = \frac{\sqrt{\mathbf{E}((CH_4^{inv.}) - (CH_4^{inv.exp}))^2)}}{\max(CH_4^{inv.exp}) - \min(CH_4^{inv.exp})} \quad (14b)$$

218 where \mathbf{E} is the mean square error function, $CO_2^{inv.}$, $CO_2^{inv.exp}$, $CH_4^{inv.}$, and $CH_4^{inv.exp}$
 219 are respectively the inventory CO_2 calculated from the numerical model, inventory CO_2 calcu-
 220 lated from experiment, inventory CH_4 calculated from the numerical model and inventory CH_4
 221 calculated from experiment. The objective function to be minimized is written as

$$\mathbf{f}(k_{irr,MGH}, k_{irr,CGH}, k_{s,CH_4}, k_{s,CO_2}) = \sum_{i=1}^4 (\text{NRMSD}_{CO_2} + \text{NRMSD}_{CH_4}) \quad (15)$$

222 where subscript i refers to experiment 1 to 4. Next using the flowchart of Babaei and Pan
 223 (2016)[Fig.4], $\mathbf{f}(k_{irr,MGH}, k_{irr,CGH}, k_{s,CH_4})$ is treated as $F(\mathbf{u}_{candidate})$. Instead of using any en-
 224 semble surrogates, we use RBF to generate surrogates of the actual solver. The number of
 225 function evaluations for Latin hypercube sampling (N_{LHS}) and the total number of function
 226 evaluations that calls the actual solver (N_{eval}) are set equal to 40 and 100. Babaei and Pan
 227 (2016) used $N_{LHS} = 2n + 2$ (where n is the number of state variables, for our case $n = 4$), and
 228 $N_{eval} = 2.5N_{LHS}$ to successfully optimize a complex problem with four variables. Therefore, in
 229 this study, $N_{LHS} = 20$ and $N_{eval} = 50$ are sufficient for optimization of objective function for
 230 four parameters (\mathbf{u}) using RBF. Furthermore, formulation of the objective function as above
 231 considers all experiments conducted in this study and both measured inventory compounds.
 232 The inventory CO_2 and methane basically include all forms of the compound in the vessel (pure
 233 (liquid or gaseous), aqueous, and hydrate), and can simply be calculated from the model as
 234 follows:

$$CO_2^{inv.} = CO_{2(liquid)} + CO_{2(aqueous)} + CO_{2(hydrate)} \quad (16a)$$

$$CH_4^{inv.} = CH_{4(gas)} + CH_{4(aqueous)} + CH_{4(hydrate)} \quad (16b)$$

235 Note that methane cannot exist in liquid form in our experimental p/T conditions.

236 3.3 Model implementation

237 The model was implemented in COMSOL Multiphysics® using its *Reaction Engineering Mod-*
238 *ule*. Two modeling setups, a batch and a reactor, were employed and coupled together. The
239 inert components (gas-hydrates and sand) were simulated by the batch model and the mobile
240 substances (water, CO₂ and methane) by the reactor model. The two modeling setups were
241 linked together to ensure a correct mass and energy balance for the entire system. The chemical
242 parameter values for hydrates and other components were taken from the literature or NIST
243 Chemistry WebBook (Linstrom and Mallard, 2013), listed in Table 1.

244 To maintain the model numerical stability, any sudden change of the boundary conditions
245 as well as shift of hydrate thermodynamics (from stable to instable and vice versa) at stability
246 temperatures must be treated continuously. To that end, the CO₂ injection intervals in model
247 were smoothed using a second derivative smoothing technique (COMSOL 4.3, 2013; Vermolen
248 et al., 2009). A rigorous method was also applied for the definition of the local reaction rates
249 (Section 3.1.3) to ensure a smooth transition of hydrate reaction rates from stable towards
250 unstable conditions.

251 The COMSOL code is converted to function $\mathbf{f}(k_{irr,MGH}, k_{irr,CGH}, k_s, CH_4$ with state variables
252 as inputs and via COMSOL-MATLAB LiveLink™, optimization is carried out in MATLAB
253 treating COMSOL as a black-box. We use MATSuMoTo toolbox in MATLAB to call RBF to
254 construct surrogate model of COMSOL function (Müller and Piché, 2011; Müller, 2014) and
255 speed up the optimization process.

256 3.4 Simulated scenarios

257 Four scenarios were simulated at the following pressure-temperature conditions where the ex-
258 perimental data are available (Deusner et al., 2012):

- 259 • experiment 1: 13 MPa/2 °C
- 260 • experiment 2: 13 MPa/8 °C
- 261 • experiment 3: 13 MPa/10 °C
- 262 • experiment 4: 8 MPa/8 °C

263 The phase diagram in Fig. 1 illustrates the experimental conditions with respect to the
264 thermodynamic stability regimes of CH₄- and CO₂-hydrate. The experiments were performed
265 in a pressure vessel of 38 cm length, 8 cm cross section diameter, 18 mm casing thickness, with
266 inlet (and outlet) of 13 mm diameter (Deusner et al., 2012).

267 The simulation developed to calibrate four experiments described above models a reactor
268 with nearly two liters volume in which 95 °C CO₂ was injected during multiple intervals separated
269 with periods of equilibrium.

270 4 Results and Discussion

271 4.1 Modeling results

272 In the present study, the major modeling results of interest are the temporal changes of (1) the
273 reactor's average temperature, (2) the overall methane and CO₂ inventory, (3) the amount of
274 methane- and CO₂- hydrate, and (4) the overall system thermal conductivity (Figs. 3 to 6).

275 In the experiments, only the total amount of inventory methane and CO₂ (including all pure,
276 dissolved or hydrate phases) was calculated using outlet and inlet volume balancing. Therefore,
277 the primary aim was to obtain a proper fit first with the methane inventory data and then
278 with the CO₂ inventory data, and then use the model to predict the fluctuations of temperature
279 and gas-hydrate in the system. Since it was very difficult to directly record temperature values
280 or determine the amount of gas-hydrates inside the pressure vessel, the use of model (after
281 constraining the unknown parameters) helped to calculate these quantities that otherwise were
282 unobtainable by means of laboratory equipments.

283 It is easy to approximately locate the start and the duration of injection intervals in Figs. 3
284 to 6 where abrupt temperature changes occurs. The system's highest temperatures are gen-
285 erally observed during the injection times when the average temperature of the system raised
286 due to the entry of 95°C CO₂. In all experiments, the majority of methane-hydrate dissociation
287 occurred during the injection intervals when the system's temperature increased above the hy-
288 drate stability temperature. Hence, the quicker the system reached or the longer it stayed at
289 hydrate instability conditions, a higher amount of hydrate dissociation was obtained. In con-
290 trast, the accumulation or precipitation of hydrate mainly occurred during equilibration periods
291 after the system lost heat to the surroundings. Further details and distinguishing features for

292 every modeling scenario are separately addressed in the following sections.

293 4.1.1 First scenario: 13 MPa/2 °C

294 The first experiment was performed at the lowest temperature leading to the lowest amount of
295 methane-hydrate dissociation and the highest amount of CO₂ accumulation. The experiment
296 time was about 44 hours in which the CO₂ was injected in four separate intervals. The maximum
297 temperature reached only 285 K and was mainly achieved at the peak of injection intervals. Due
298 to very low ambient temperature and high vessel pressure, the system hydrates were exposed to
299 instability conditions only for a very short time. Most of the CO₂ was, therefore, speculated to
300 deposit in the vessel as CO₂-hydrate with excess pore water. The modeling results also confirmed
301 the accumulation of CO₂-hydrate in the system. The qualitative model reproduction of the
302 experiment data of the CO₂ inventory supports this hypothesis and also suggests a homogeneous
303 retention of the injected CO₂ in the vessel (Fig. 3).

304 The long equilibration periods between the injection intervals allowed CO₂ to slowly form
305 CO₂-hydrate and increased its retention yield. **The model predicted the formation of nearly**
306 **3 mol CO₂-hydrate inside the vessel while the methane-hydrate dissociation was predicted to**
307 **be less than 0.1 mol.** A substantial formation of CO₂-hydrate with the excess pore water was
308 confirmed and was speculated as the main reason preventing rapid growth of preferential flow
309 paths in this scenario.

310 4.1.2 Second scenario: 13 MPa/8 °C

311 In comparison to the first experiment, the second experiment was done at a higher ambient
312 temperature and therefore a significantly higher amount of methane-hydrate dissociation was
313 observed (Fig. 4). While the length of experiment was marginally longer than the first experi-
314 ment (about 45.5 hours), a higher amount of CO₂ was injected through five intervals (25% more
315 CO₂ was injected in comparison to the first experiment). The amount of heat transferred to the
316 vessel was therefore higher but this was not the only feature contributing to a higher amount
317 of methane-hydrate dissociation. In this scenario, the system was exposed to the hydrate insta-
318 bility conditions for a longer time thereby increasing the methane yield. Evidently, the small
319 temperature difference between the experiment's initial condition and the hydrates instability
320 zone derived the system to gas-hydrates instability conditions faster and led to the dissociation

321 of a larger amount of methane-hydrate (the second highest amongst all experiments). The CO₂
322 inventory was overestimated by the model. This suggests that the injected CO₂ was possibly
323 conveyed through preferential flow paths that were created due to methane-hydrate dissocia-
324 tion. Other factors such as consecutive injections, and a short equilibration time between the
325 injection intervals, could also have enhanced the progression of the preferential flow paths in
326 this experiment.

327 4.1.3 Third scenario: 13 MPa/10 °C

328 The third experiment was done at the highest ambient temperature during which the CO₂-
329 hydrate was subjected to instability conditions for the entire duration of the experiment and
330 therefore never formed. The modeled CO₂ inventory curves deviated even more from the ex-
331 perimental data indicating once again the development of preferential flow paths prohibiting a
332 spread of CO₂ into the reactor volume. Comparing the results with the previous scenarios, it
333 is speculated that the development of preferential flow paths are even stronger and that the
334 formation of such pathways can be a function of ambient system temperature. Modeling results
335 predicted a slightly higher dissociation of methane-hydrate than the second experiment while in
336 the reality it was lower (Fig. 5). It is speculated that in the absence of CO₂-hydrate formation,
337 the injected CO₂ at later stages followed the formerly generated pathways and discharged faster
338 from the outlet. However, this was not the case for the second scenario where the slight forma-
339 tion of CO₂-hydrate during the equilibrium intervals might have plugged the previously formed
340 pathways, forcing the injected CO₂ in the following stages to spread into the regions with high
341 methane-hydrate concentration.

342 The modeling of this scenario revealed that the strongly developed and hydraulically con-
343 nected preferential flow paths dramatically disturbed the uniform distribution of the heat that
344 was introduced via the injection of supercritical CO₂. Therefore, the interactions between the
345 injected CO₂ with the remaining methane-hydrate in the vessel was limited. Most of the heat at
346 later injection intervals was, thus, expelled from the system and, despite the higher experimental
347 ambient temperature, a lower rate of methane-hydrate dissociation was achieved.

348 The results show a clear dissimilarity between modeled and experimental data since the
349 beginning and particularly after the consecutive first and second injection intervals. The quick
350 progression of the preferential paths in this scenario may thus not only be related to the absence

351 of CO₂-hydrate formation but also may be favored by the consecutive injections at the beginning
352 of the experiment. The total duration of this experiment was about 77 hours, the longest amongst
353 all.

354 4.1.4 Forth scenario: 8 MPa/8 °C

355 The fourth experiment (Fig. 6) was performed at a lower pressure compared to previous experi-
356 ments. The ambient temperature as shown in Fig. 1 was slightly below the stability temperatures
357 of both CO₂- and methane-hydrate and equal to that in the second scenario. The system thus
358 easily reached hydrate instability conditions during the CO₂ injection intervals. The highest
359 amount of methane dissociation was achieved in this experiment given its total duration was
360 longer (about 50% longer) than the second experiment. The formation of preferential flow paths
361 is evident as a result of CO₂ inventory mismatch. The quick progress of preferential flow paths
362 after the consecutive injections of CO₂ at the second and third intervals is visible. As for the
363 second experiment the formation of CO₂-hydrate favored the distribution of the injected CO₂
364 and enhanced the overall methane-hydrate dissociation in comparison with the third experiment.
365 The experiment took roughly 71 hours to complete.

366 4.2 Estimated kinetic parameters

367 Most of the parameter values were taken from the literature, or calculated by SUGAR toolbox
368 (Kossel et al., 2015) in close approximation with the previously reported values (see Table 1).
369 The only unknown parameters that often vary between different systems were $k_{irr,MGH}$, $k_{irr,CGH}$
370 , k_{s,CH_4} , and k_{s,CO_2} . Using the above described optimization technique we obtained the fol-
371 lowing values for these parameters $k_{irr,MGH} = 0.02(\text{mol.m}^{-3}\text{s}^{-1})$, $k_{irr,CGH} = 0.03(\text{mol.m}^{-3}\text{s}^{-1})$,
372 $k_{s,CH_4} = 4 \times 10^{-5}(\text{s}^{-1})$, and $k_{s,CO_2} = 1 \times 10^{-5}(\text{s}^{-1})$. This values are in agreement with previously
373 reported values. For instance, the values of k_{s,CH_4} and k_{s,CO_2} are in the same range of values
374 reported in Noyes et al. (1996) for first-order gas-exchange rate constant. The estimated values
375 for parameters describing hydrates dissociation at absolute instability conditions, $k_{irr,MGH}$ and
376 $k_{irr,CGH}$, were about two orders of magnitude lower than the value reported in Jerbi et al. (2010)
377 for CO₂ dissociation. However, Jerbi et al. (2010) performed the experiments in a semi-batch
378 stirred tank reactor at stirring velocity of 450 rpm. A simple comparison between the two sys-
379 tems (pressure vessel and stirred-tank reactor) shows that it is reasonable to obtain significantly

380 lower dissociation rates in a pressure vessel.

381 We were able to obtain a convenient fit to methane inventory data for all the scenarios except
382 the third one performed at 13 MPa/10 °C. The fact that neither the model results for methane
383 inventory nor the results for CO₂ inventory of the third scenario were found to reasonably fit to
384 the experimental data (Fig. 5) suggests that the underlying processes in this experiment were
385 too complicated to be described by the modeling approach presented. It is therefore difficult
386 from this approach to correlate the rate of methane-hydrate dissociation in the third experiment
387 with those in other scenarios. On the contrary, the model results did not fit properly to the CO₂
388 inventory data at all. This might be mainly due to the development of preferential flow paths
389 in the system causing the CO₂ to poorly spread in the reactor volume and to leave the reactor
390 early. Since the current model assumptions are based on perfect mixing, any deviation of model
391 results from the CO₂ inventory data can be linked to the occurrence of preferential flow paths
392 and the heterogeneous transport of CO₂ inside the vessel.

393 The aim was not, however, to obtain a perfect fit to each experiment with any combination of
394 the values, but to find for each of these parameters a constant value to which a reasonable fit can
395 be achieved to all the scenarios. It is worth noting that most of the parameters in reality might
396 be a function of temperature, pressure and salinity. Since pressure was kept nearly constant in
397 the vessel and the temperature of the system fluctuated within a narrow band, the majority of
398 chemical properties including the estimated effective rates were assumed constant.

399 4.3 Dissolution rate of methane and carbon dioxide in water

400 A significant sensitivity of the model to the dissolution rates of methane and CO₂ in water was
401 found during the model analysis. It was displayed that not only the final aqueous concentrations,
402 c_{cte,CH_4} and c_{cte,CO_2} (calculated from SUGAR toolbox (Kossel et al., 2015) and listed in Table 1),
403 but also the exchange rate constants between water and gas, k_{s,CH_4} and k_{s,CO_2} , are important
404 for the dissociation/formation of the hydrate at the beginning of the experiments and after
405 the injection intervals. Numerical stability of the model was found to be very sensitive to the
406 values of these parameters. These parameters might be less influential at field-scale than in
407 the experiments due to the comparatively larger computation time-scales or larger size of the
408 domain.

409 4.4 The relation between carbon dioxide retention and methane release

410 The present findings speculate that the differences between modeling and experimental results
411 are associated with the presence of flow heterogeneities and their relative growth inside the
412 vessel. Since model predictions are based on perfect mixing assumptions, the differences between
413 model predictions and experimental data in CO₂ inventory curves (Figs. 4 and 6) indicate that
414 the injected CO₂ bypassed the majority of the vessel contents in all scenarios except the first
415 one. This, however, only hindered the methane-hydrate dissociations in the third experiment,
416 suggesting that in both, the second and fourth scenarios the injected CO₂ still managed to
417 deliver its heat to the vessel contents.

418 The only major difference between other scenarios and the third scenario is the formation
419 of CO₂-hydrate, which appears to affect positively the dissociation of methane-hydrate in the
420 second and the fourth scenarios. The reason may be related to the formation of solidified CO₂-
421 hydrate clogging up the previously formed flow pathways, thereby forcing the upcoming CO₂
422 to choose a different pathway. Alternatively, the lower enthalpy of formation of CO₂-hydrate in
423 comparison to methane-hydrate may have thermodynamically favored methane-hydrate disso-
424 ciation. Either way, it appears that at p/T conditions closer to the methane-hydrate instability
425 zone, the retention of CO₂ catalyzes the methane-hydrate dissociation. This might have the
426 following practical implications for CO₂ injection into hydrate reservoirs. First, the tempera-
427 ture of the injected CO₂ can be adjusted in order to avoid the reservoir temperatures at which
428 CH₄-hydrate is stable and CO₂-hydrate is unstable (at p/T conditions similar to the third sce-
429 nario). Secondly, altering the reservoir conditions to the p/T conditions at which CH₄-hydrate
430 is unstable and CO₂-hydrate is stable might increase both CO₂ retention yield and CH₄-hydrate
431 dissociation. This might be only obtainable by combining the two techniques of depressurization
432 and thermal stimulation together. Accordingly, it might be safe to say that injection of CO₂ into
433 the gas-hydrate reservoirs at p/T conditions similar to the third scenario is not economically
434 and environmentally favorable.

435 4.5 Guidelines for field-scale modeling

436 The dissociation of hydrate was not entirely related to the amount of heat that was introduced
437 to the system but to the quality of heat distribution, information that is difficult to quantify em-
438 pirically. The comparison between an ideally mixed model and experimental data allowed us to

439 interpret the system behavior in each case based on the discrepancies observed. Methane-hydrate
440 dissociation yield is likely related to the relative combination of several factors that cannot be
441 imposed externally, such as reservoir temperature, pressure, salinity, structural heterogeneities,
442 composition of sand layers, or the spatial distribution of these quantities. In turn, several factors
443 can be regulated during the production from gas-hydrate deposits which were discovered to have
444 a noteworthy influence on the final results. These include temperature of superheated CO₂, the
445 equilibration periods between injection intervals, and the injection strategy. The succession of
446 injection intervals was shown detrimental to the whole process due to bolstering the preferential
447 flow paths and thus decreasing the quality of heat expansion. The mismatches between model
448 and experiments were mostly observed after the consecutive injections. The longer the equili-
449 bration intervals, the lesser was the extension of preferential pathways through the vessel. Low
450 injection rates of CO₂ were found to benefit the CO₂ retention process through homogenizing
451 the distribution of CO₂, allowing it to disperse further into the depth of hydrate deposit while
452 preventing the restoration of preferential flow paths. A similar finding has been recommended
453 by [White et al. \(2011\)](#). It is also suspected that the formation of CO₂-hydrate not only im-
454 proved the quality of CO₂ retention but also enhanced the overall methane-hydrate dissociation.
455 Therefore, the method at p/T conditions between the two hydrates stability curves (at condi-
456 tions similar to the third experiment) was shown highly ineffective. However, more data are
457 needed to prove that the CO₂-CH₄-hydrate conversion must be avoided at such p/T conditions
458 by performing more experiments at such conditions.

459 In addition, it was shown that the pore-scale heterogeneities that are typically ignored at
460 field-scale models can immensely affect the simulation results. Since the inclusion of such small-
461 scales effects into reservoir models is computationally very elaborate, the urge of upscaled models
462 which are able to estimate the small-scale (cm to meters) dynamics in the presence of hetero-
463 geneities as a function of system pressure and temperature is noted. These models can be either
464 empirical or analytic.

465 4.6 Model predictability and limitations

466 It was shown that the model performs better at low temperatures and high pressures deep
467 inside the hydrate stability zone (at conditions similar to the first experiment). However, the
468 predictability of model reduced at higher temperatures, closer to hydrate instability zone (at

469 conditions similar to the third experiment). The analysis of results shows that the presented
470 model was able to forecast the behavior of a (semi-)homogenized system. A similar findings
471 was noted by the Ignik Sikumi field trial (Schoderbek et al., 2013). The deviations between
472 modeling results and experiments occurred when preferential flow paths played a major role on
473 the transport of substances, and when the system become extremely heterogeneous.

474 5 Conclusions and Implications

475 We presented here the results of a kinetically-focused simulation which is used to explain the ex-
476 perimental results reported in Deusner et al. (2012) without extra complexities of fully spatially-
477 resolved, computationally-expensive fluid dynamics simulations. Unlike most of the studies in
478 this field in which the focus has been given to the fluid dynamics and transport effects and as
479 a result reaction kinetics were oversimplified, in this study a detailed definition of kinetics was
480 employed. The transport phenomena, however, were simplified to a basic mass and energy bal-
481 ance for a representative elementary volume (2 liters) which is equal or smaller than the typical
482 size of a grid block in continuum field-scale models. This study demonstrates the significance of
483 reaction kinetics on the extraction of natural gas through the injection and exchange of CO₂ with
484 methane in gas-hydrates. The details of kinetics are therefore shown to be too significant to be
485 easily discarded despite the fact that such simplifications are commonly observed in field-scale
486 models. Furthermore, it is noted that an equal emphasis should be given to the details of small-
487 scale heterogeneities in reservoir simulators in order to correctly model hydrate exploitation at
488 field-scale. To avoid excessive computational demands of high-resolution models at field-scale
489 while taking the effects of pore-scale heterogeneities into consideration, it is required to develop
490 upscaled models that perform at small scales (cm to m). Such up-scaled models do not ex-
491 plicitly solve all the transport mechanisms in details, but describe and encapsulate the overall
492 impact of small-scale heterogeneities into a relatively accurate and computationally-inexpensive
493 box-model. With the help of the model we estimated the values of key intrinsic parameters that
494 are unknown, and are different depending on the experimental setup employed. These parame-
495 ters are usually difficult to directly quantify from the experiments and as a result often over or
496 underestimated.

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601 **Tables**

Table 1: Parameter values used in this study. Some of the parameter values are recalculated using SUGAR toolbox (Kossel et al., 2015) for the p/T conditions and the salinity of each scenario. For the rest of parameters the values from literature are taken as default.

Parameter (Unit dimension)	Scenario 1	Scenario 2	Scenario 3	Scenario 4	Literature values
Heat capacity of CH ₄ -hydrate ($Jmol^{-1}K^{-1}$)	251.18	268	275.8	268	258 ^a
Heat capacity of water ($Jmol^{-1}K^{-1}$)	73.8	73.91	73.95	74.16	75.4 ^b
Heat capacity of CO ₂ -hydrate ($Jmol^{-1}K^{-1}$)	307.77	328.38	337.95	328.38	
Heat capacity of methane ($Jmol^{-1}K^{-1}$)	59.5	57.21	56.53	47.69	34.8 ^c
Heat capacity of CO ₂ ($Jmol^{-1}K^{-1}$)	92.97	95.71	96.77	106.62	103 ^d
Heat capacity of sand (Quartz) ($Jkg^{-1}K^{-1}$)					730 ^d
Heat transfer coefficient of the steel jacket ($Wm^{-1}K^{-1}$)					16 ^e
Thermal conductivity of water ($Wm^{-1}K^{-1}$)	.575	.585	.588	.583	0.58 ^f
Thermal conductivity of CO ₂ ($Wm^{-1}K^{-1}$)	.123	.116	.114	.108	0.1 ^g
Thermal conductivity of CH ₄ ($Wm^{-1}K^{-1}$)	.05	.05	.05	.042	0.051 ^h
Thermal conductivity of CH ₄ -hydrate ($Wm^{-1}K^{-1}$)					0.62 ^h
Thermal conductivity of CO ₂ -hydrate ($Wm^{-1}K^{-1}$)					0.49 ^d
Thermal conductivity of water-saturated sand ($Wm^{-1}K^{-1}$)					3.3 ⁱ
CH ₄ -hydrate enthalpy of dissociation ($kJmol^{-1}$)					54.2 ^o
CO ₂ -hydrate enthalpy of dissociation ($kJmol^{-1}$)					82 ^j
Activation energy of CH ₄ -hydrate dissociation ($kJmol^{-1}$)					81 ^j
Activation energy of CH ₄ -hydrate formation ($kJmol^{-1}$)					20.6 ^k
Activation energy of CO ₂ -hydrate dissociation ($kJmol^{-1}$)					102 ^j
Activation energy of CO ₂ -hydrate formation ($kJmol^{-1}$)					20 ^l
Dissolution kinetic constant of CH ₄ -hydrate ($mol\ m^{-2}\ s^{-1}\ Pa^{-1}$)					3.6×10^4 ^m
Formation kinetic constant of CH ₄ -hydrate ($mol\ m^{-2}\ s^{-1}\ Pa^{-1}$)					0.6 ⁿ
Dissolution kinetic constant of CO ₂ -hydrate ($mol\ m^{-2}\ s^{-1}\ Pa^{-1}$)					1.83×10^8 ^j
Formation kinetic constant of CO ₂ -hydrate ($mol\ m^{-2}\ s^{-1}\ Pa^{-1}$)					5×10^{-3} ^o
Equilibrium concentration of CH ₄ , $c_{cte,CH4}$ ($mol\ m^{-3}$)	156.36	139.34	134.5	102.6	
Equilibrium concentration of CH ₄ -hydrate, $c_{cte,MGH}$ ($mol\ m^{-3}$)	58.49	84.33	95.53	85.51	
Equilibrium concentration of CO ₂ , $c_{cte,CO2}$ ($mol\ m^{-3}$)	1884.05	1695.93	1642.09	1625.52	
Equilibrium concentration of CO ₂ -hydrate, $c_{cte,CGH}$ ($mol\ m^{-3}$)	909.1	1376.34	1617.89	1386.51	

^aHanda (1986)

^bKoh et al. (2011)

^cWaite et al. (2009) at 260K

^dJung et al. (2010)

^eVessel manual

^fWaite et al. (2009)

^gNIST Chemistry WebBook (Linstrom and Mallard, 2013)

^hWaite et al. (2007)

ⁱTarnawski et al. (2011)

^jClarke and Bishnoi (2004)

^kFreer et al. (2001)

^lEstimated

^mClarke and Bishnoi (2001); Kim et al. (1987)

ⁿEnglezos et al. (1987)

^oClarke and Bishnoi (2005)

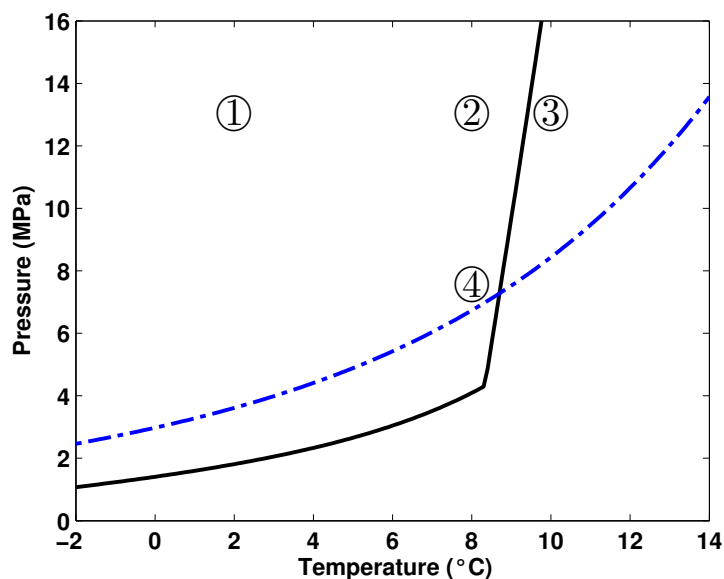


Figure 1: Stability curves for methane-hydrate (blue dotted line) and CO_2 -hydrate (black solid line) in respect to system pressure and temperature. At high pressures and low temperatures inside the stability zone hydrate dissolves/forms according to Reactions 8a and 8b. Outside the hydrate stability zones at high temperatures and low pressures hydrate only dissociates (Reactions 9a and 9b). The encircled numbers indicate the pressure and temperature conditions at which the experiments were performed.

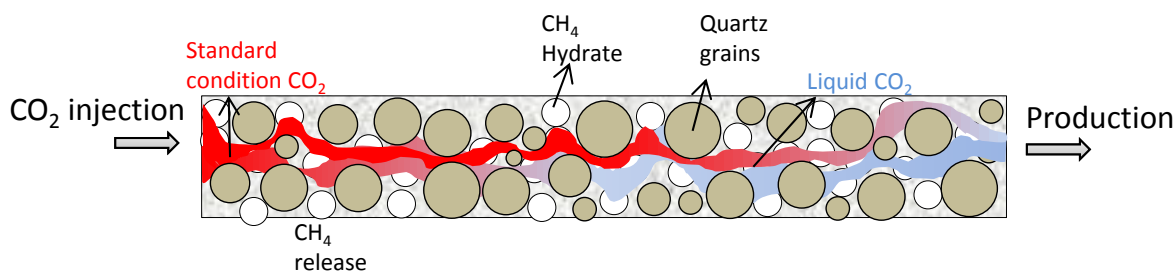


Figure 2: The schematic diagram of the processes occurring during supercritical CO_2 injection into the pressure vessel containing CH_4 -hydrate and sand.

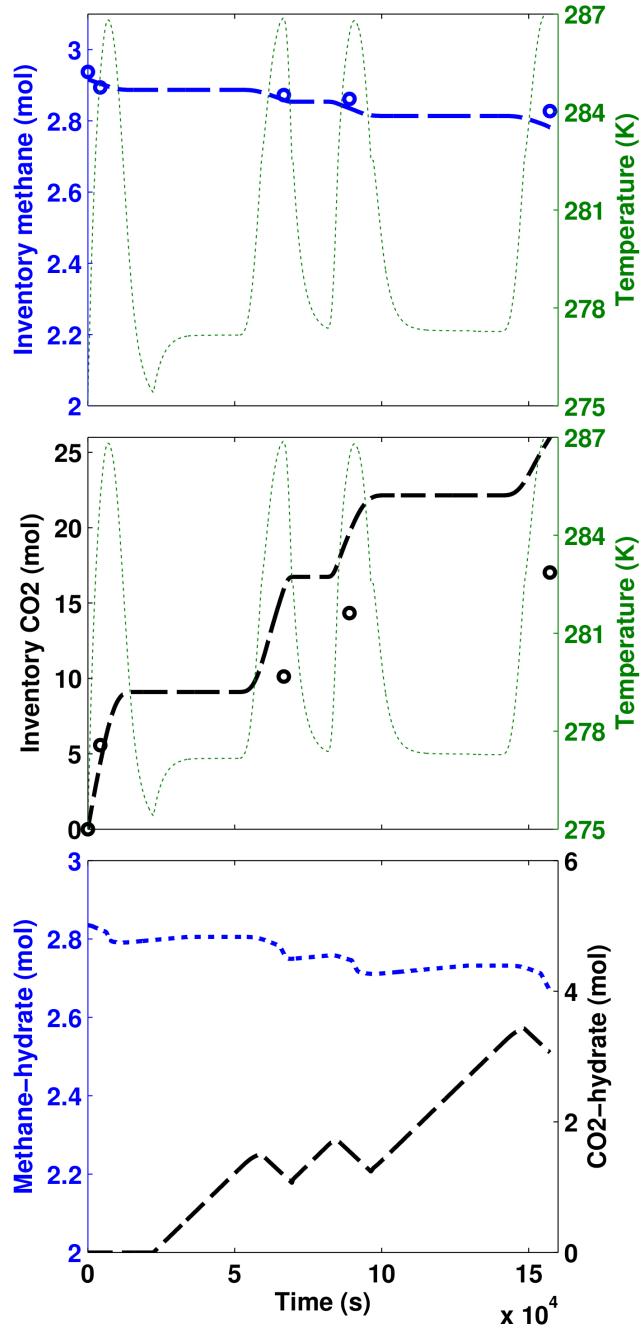


Figure 3: The modeling results of first scenario are shown by dashed lines and open dots are the reported experimental data for the p/T conditions of 13MPa and 2°C.

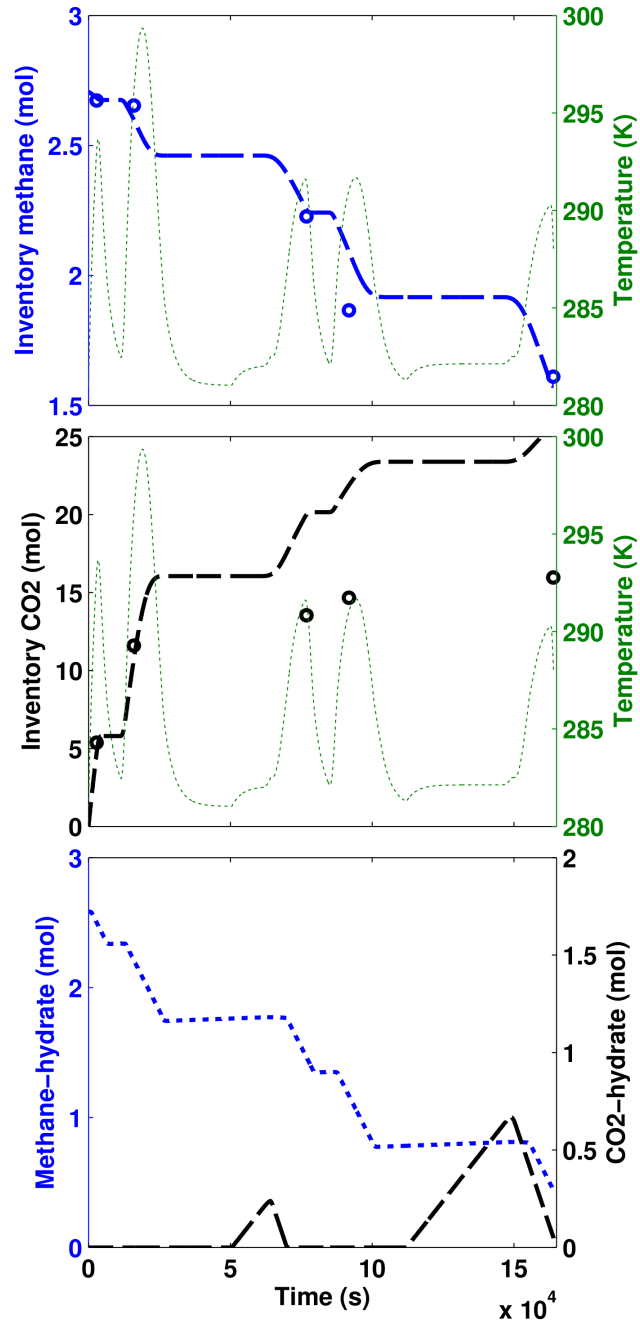


Figure 4: The modeling results of second scenario are shown by dashed lines and open dots are the reported experimental data for the p/T conditions of 13MPa and 8°C.

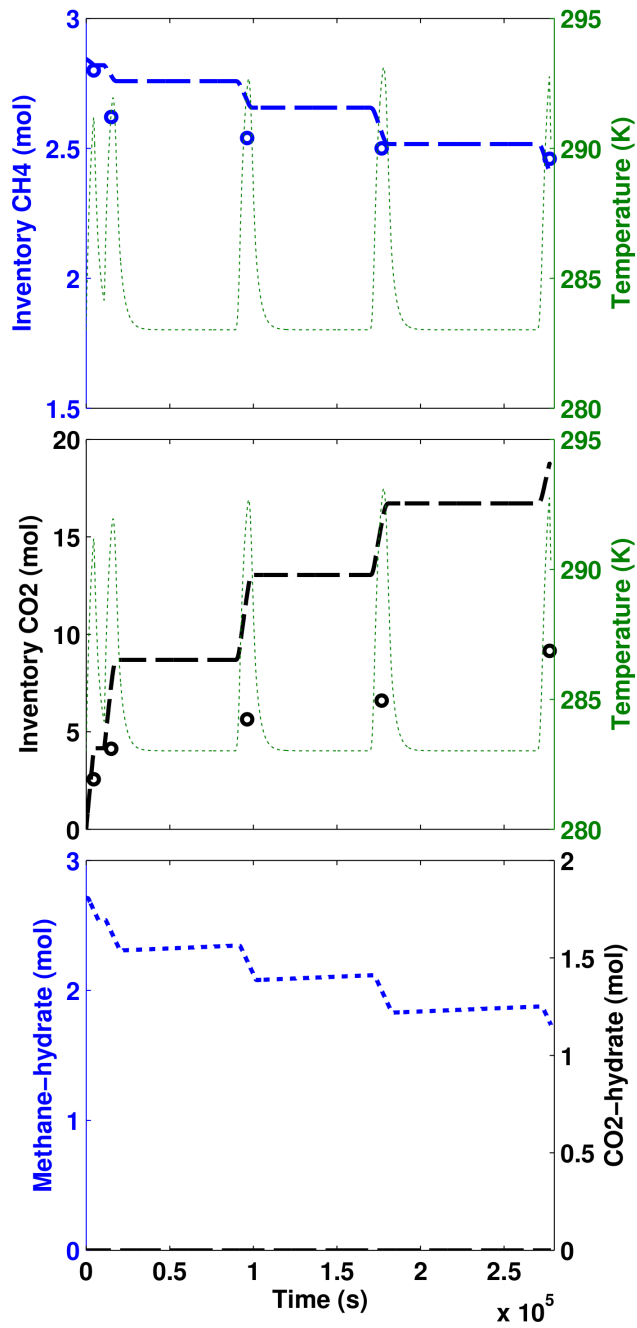


Figure 5: The modeling results of third scenario are shown by dashed lines and open dots are the reported experimental data for the p/T conditions of 13MPa and 10°C.

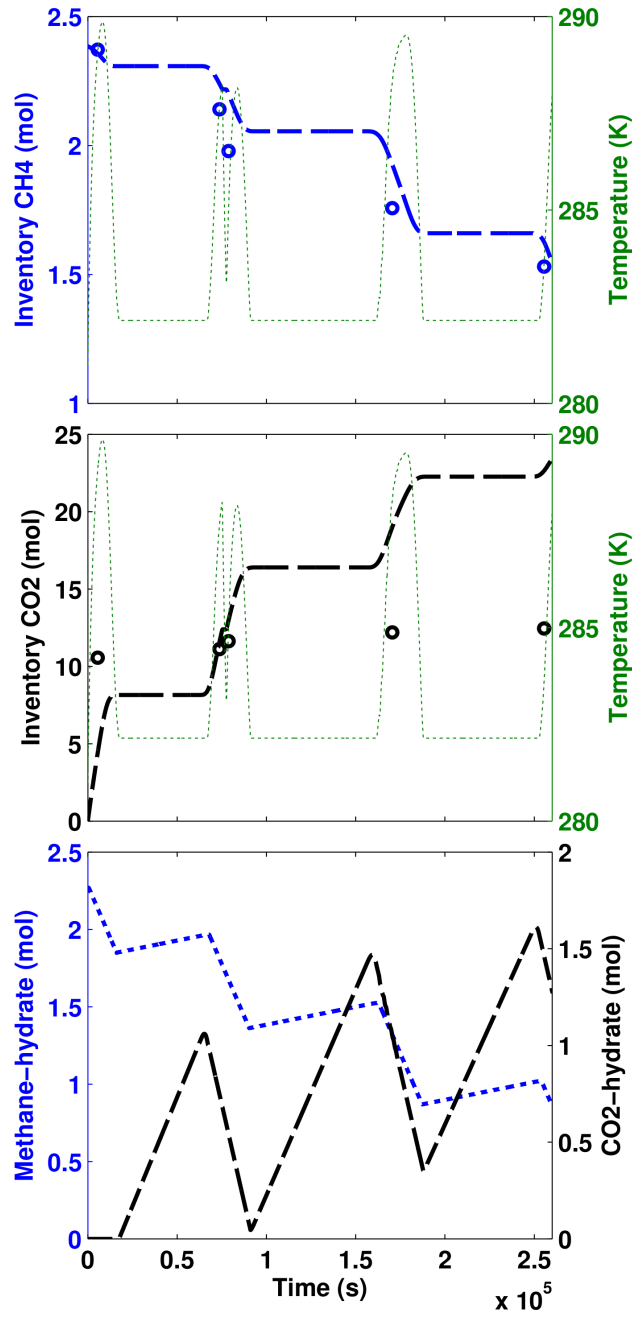


Figure 6: The modeling results of forth scenario are shown by dashed lines and open dots are the reported experimental data for the p/T conditions of 8MPa and 8°C.