

Melt-induced buoyancy may explain the elevated rift-rapid sag paradox during breakup of continental plates

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Timing of breakup relative to the sag and salt intervals in central South Atlantic

Disparities in the rates and timing of subsidence at rifted margins from those predicted by McKenzie's rift-cooling model¹ are expressed as thinner than expected syn-rift thicknesses, greater than expected early post-rift thicknesses and subaerial-shallow water conditions at breakup rather than deep water²⁻²⁶. These issues are well known in the central South Atlantic region where, for example, the early post-rift sag and salt section is >4 km thick in northern Kwanza and composed of shallow water-emergent carbonates and evaporites. This has led some researchers to argue that the sag and salt intervals (Aptian, 123-113 Ma) are actually late syn-rift^{3,15,27}. However, neither water depths nor basin geometry really support the idea (e.g. Model tR113, Figs.S1 and S2). In particular, the absence of extensional growth faults in the sag^{2-5,10,15,27} (Fig. S3) and the geometric requirement of large trans-continental shears in plate reconstructions to accommodate late breakup^{24,25,28}, leads us to strongly favour breakup at c.123 Ma (pre-sag). This interpretation is supported by reported ages of syn-rift, sag and mid-ocean ridge-type basalts in wells drilled at or close to the point of breakup in deep water Angola by Statoil, ConocoPhillips and Total in 2014-2015 (e.g. Dilolo-1, Jacare-1, Kamoxi-1, Umbundu-1) and by more recent plate reconstructions^{5,29,30} (Movie S4).

Information on the ages of rifting

The best estimate of the onset of rifting is 131 Ma based on the oldest syn-rift sediments which are Hauterivian in age¹⁶. The precursor Paraná volcanics are precisely constrained by ⁴⁰Ar/³⁹Ar dating to 135 Ma +/- 1 My³¹ and uniform layering within them shows that they were erupted before extension began. Information on the timing of the end of rifting (123 Ma) is reported in the main article and in Refs.^{2,5,8,16,18,30}.

Seismic interpretation

The nature of the seismic units in the central South Atlantic region has previously been described by Ref.⁵. The syn-rift interval is characterized by tilted and fanning reflection geometries. Strong parallel reflections at the base of the syn-rift interval correlate with the top of basalt (the major Paraná eruptive event, the precursor to rifting in the central South Atlantic region^{8,31}). The top of the syn-rift interval is marked by an angular unconformity onto which flat-lying sag strata progressively overstep (Fig. S3). The top of the sag is a flat horizon above which the salt has a clear but chaotic seismic expression resulting from salt tectonic movement. The top of the salt is marked by a strong seismic reflection and is overlain by Albian carbonates characterized by growth, fault and fold geometries. These structures were formed during gravity spreading as the salt tilted towards the ocean as a result of enhanced subsidence in that direction¹⁸.

Are anomalous subsidence patterns during breakup related to mantle plumes?

The aim of our work is to find a generic solution for the common problem of elevated breakup and subsequent collapse at rifted margins^{6,11}. We acknowledge that the thermal and dynamic effect of a mantle plume can have a semi-regional effect³²⁻³⁴. However, we feel that this is not a universal explanation for the subsidence issue associated with McKenzie's model¹ during the breakup of continental plates. Nonetheless, one of us (D.G.Q., Ref.⁵) has previously promoted this idea for a region of the central South Atlantic lying south of the current study area (Fig.S3), where syn-rift strata are particularly thin. It is therefore feasible that anomalous uplift caused by mantle plumes may enhance buoyancy caused by melt and phase changes in the mantle but probably not along every rift margin.

Information on melt produced during rifting

In the accompanying article, we explore the idea that melt is retained in the asthenosphere during rifting and lost rapidly after breakup. That melt exists in the mantle beneath modern continental and oceanic rifts is recognized from an increase in electrical conductivity, a reduction in seismic velocities and a decrease in gravity³⁵⁻⁶⁰. The geophysical evidence for melt is also supported by theoretical models, experiments and geochemical data⁶¹⁻⁷³. Unfortunately, the amount of melt is difficult to constrain because many properties of the asthenosphere such as interconnectivity of melt-filled pores and permeability are poorly constrained with large uncertainties in physical characteristics, including the shape and alignment of melt inclusions^{40,43,74-77} and whether the peridotite is in textural equilibrium⁷⁸⁻⁸⁰.

As geophysical data and geodynamic models grow, two different hypotheses have emerged on the amount of melt present beneath modern rifts, with continental and oceanic rifts often treated similarly.

- 1) A common view is that dihedral angles of melt formed in olivine-rich peridotite are significantly less than 60°^{79,81} and the peridotite is texturally equilibrated⁷⁸ meaning that melt flows through the matrix at very low porosities. Consequently, the mantle will retain less than c.2% melt^{35,62,70,83-86}, migrating out of the source region faster than it is generated^{67,87}. This view is supported by various geophysical interpretations of <1%-3% melt in regions of upwelling asthenosphere⁴³, more commonly in oceanic domains^{48,50,54-56,59,88}.

- 2) Alternative experimental work and theoretical models indicate that lherzolite can actually retain 5%-c.20% melt^{79,80,89,90}, consistent with other geophysical interpretations of c.5-16% melt in zones 30-100 km thick and 30->150 km wide in the upper mantle^{44,46,47,91}, including continental regions⁵². For example, Ref.⁹², using magnetotellurics, and Ref.⁷⁷, using P-S wave velocities and anisotropy, have independently produced similar estimates of 6-15% melt over an interval 30-60 km thick in the upper mantle beneath the Dabbahu rift segment in Afar, Ethiopia, equivalent to a melt column of 3-10 km. Assuming these interpretations are valid, there are five potential reasons for high melt retention.
 - a) Where peridotite is neither mono-mineralic nor equi-granular, melt does not become interconnected until porosities of somewhere between 5-20%^{79,80,90}.
 - b) Textural equilibrium is difficult to establish in flowing mantle⁷⁷.
 - c) Interfacial tension presents a significant barrier to compaction- and buoyancy-driven flow at porosities of less than c.20%, even where grain boundaries are wet^{89,93}.
 - d) Viscous resistance to flow means migration can be slower than generation^{86,94}.
 - e) Melt flowing between crystals or carried upwards in circulating asthenosphere is trapped by a viscous lid of mantle lithosphere which represents a permeability barrier^{39,56,88,95,96}, ultimately resulting in crystal fractionation⁹⁶⁻⁹⁸.

One of the key points is that melting experiments on real peridotite show there is a drastic reduction in the interconnectivity of pores when minerals other than olivine are present⁸⁰. As much as 21-29% basaltic melt can be retained in lherzolite containing less than c.80% olivine⁹⁰ up to the point of disaggregation⁸⁹. A typical fertile lherzolite comprises 60% olivine, 20% orthopyroxene, 10% clinopyroxene and 10% garnet or spinel or plagioclase⁹⁹, so theoretically all the aluminium-rich phases, all the clinopyroxene and some orthopyroxene have to melt before the pores become fully connected. Thus, it may be that melt is retained in asthenosphere composed of heterogeneous lherzolite but will tend to flow out when it is mostly olivine, provided interfacial tension is overcome⁸⁹. However, in the continental domain a more significant permeability barrier exists at the top of the asthenosphere in the form of the continental plate^{95,96,98}. Magma can only ascend through the overlying lithospheric mantle in discrete, fracture-assisted intrusions rather than by porous flow^{9,43,74,84}. These considerations offer a way of aligning models of low versus high melt retention beneath rifted margins.

- Prior to breakup, when rifting reaches a peak, large amounts of melt have been generated and a significant proportion is temporarily retained in the upwelling asthenosphere, initially because porous flow is inhibited in multi-mineralic, non-equilibrated lherzolite and thereafter because the overlying lithosphere forms a closed lid. The zone of melt-impregnated asthenosphere will tend to drain downwards by capillary action, a wetting profile driven by grain-grain tension.
- After breakup, the mantle down to c.90 km below the inactive rifted margins is largely devoid of garnet-spinel-plagioclase and clinopyroxene, these minerals having already been lost to melt. The depleted asthenosphere consists predominantly of olivine-rich, texturally-equilibrated harzburgite-dunite with a dihedral angle significantly less than 60° meaning most of the remaining melt is interconnected at low porosities (<1-3%) and migrates relatively quickly to the base of the lithosphere. Provided the melt does not solidify, it will then flow up-dip towards the edge of the continental plate to be incorporated in thick oceanic crust created at the new spreading axis. As the depleted asthenosphere cools, it is gradually accreted to the base of the continental plate to form new lithosphere together with melt which has solidified before reaching the oceanic domain.

With these processes in mind, we suggest it is possible that asthenosphere retains 5-16% melt during the latter stages of continental rifting prior to breakup^{52,77,92}, more than below mid-ocean ridges where there is no mantle lid. In our melt models, we actually use average amounts at the lower end of this range (those prefixed M and F in Figs. 1 and 2 in the accompanying article, Figs.S1 and S2 and Movies S1-S3).

It is also worth noting that the presence of melt beneath an extending plate is likely to cause periodic intrusions of sills or underplates into the lithosphere⁹⁷, an effect which may partly compensate for thinning of the plate, allowing certain rifted margins to become surprisingly wide (e.g. Fig.8 in ref.⁵). Underplates are quite well documented beneath central South Atlantic rifted margins with interpreted thicknesses up to 6 km^{14,25}.

Based on the arguments above, we test the effect of relatively high melt retention in the asthenosphere. The idea is that beneath continental rifts up to breakup i) significant porosity is required for melt to become interconnected in the asthenosphere^{76,79,80,89,90} and ii) the base of the continental lithosphere acts as a lid^{39,57,88,95,96,98}.

Once it becomes interconnected, melt migrates upwards through the asthenosphere by porous flow, controlled by compaction and buoyancy, viscosity, permeability and surface tensions^{62,67,89,93,94,99}. The upward migrating melt will be impeded by continental mantle lithosphere and the zone of melt-impregnated asthenosphere will then increase in thickness downwards by capillary action, with wetting driven by grain–grain tension to give a smoothly decreasing melt profile^{89,93}, as used in our models (Fig. 1c in main article).

We choose to limit the average amount of melt in the asthenosphere to 8%, starting with 0% at 90 km depth up to a maximum of 16% at 20 km close to the point of breakup where the lithosphere is highly attenuated (stretching factor of >7.5, Figs. S4c and S4d). This is equivalent to a maximum column of 5.6 km melt. However, over most of the Kwanza Basin, where stretching factors are <5 (Fig. S4c), a column less than half this amount has developed in our models by the end of rifting. For comparison, more melt is involved in the construction of oceanic crust but here it ascends relatively quickly to a mid-ocean ridge, unimpeded by lithosphere. A total column of 7.5 km melt is needed for average oceanic crust^{73,82,97} but as much as 12-24 km of melt may be involved during the formation of volcanic rifted margins^{9,87,96}.

We do not argue for unusual mantle conditions of temperature, only that a significant proportion of melt cannot migrate out of the asthenosphere until after breakup. In our models, the asthenosphere wells up passively while the lithosphere extends, a general assumption in most models of continental rifts^{13,83}. In reality, the presence of melt in the asthenosphere may enhance convection due to buoyancy and a possible reduction in viscosity^{60,61,83,99,100} although Ref.¹⁰¹ suggests that viscosity may initially increase as the asthenosphere becomes drier with melting. If active upwelling or a mantle plume are associated with rifting, then higher degrees of melt are expected due to replenishment of fertile peridotite⁹⁶ or higher temperatures^{64,82} which will lead to more volcanism at breakup⁸⁷.

Information on loss of melt after breakup

During rifting, significant columns of melt may be trapped beneath continental lithosphere until it breaks, after which it is likely to drain rapidly to the new volcanic ridge^{87,94,102}. This may explain the release of huge volumes of magma over short periods of time forming thick crust on volcanic margins^{82,96,97,102} as well as outer volcanic highs on non-volcanic margins^{5,23,103}. In the North Atlantic Igneous Province, Ref.¹⁰⁴ suggests that a c.7 km column of melt per unit area was released to the crust in the 5 My up to breakup followed by 6-17 km soon after breakup, which we suggest was partly due to lateral migration of melt stored in the asthenosphere (cf. Ref.¹⁰⁵).

After breakup, the asthenosphere beneath the rifted margins ceases to flow and textural equilibrium is re-established. Melt production also stops and porosity decreases causing the rate of melt migration to decline. The density of the mantle will increase as melt is lost, causing the rifted margins to subside. Based on empirical observations of post-breakup volcanism and thicknesses of early post-rift sag⁵, we expect melt to be lost at an exponential rate with only small percentages left after 10 My, although the lithosphere continues to thicken and subside by cooling.

Basin reconstruction approach

TecMod is modeling software for lithosphere-scale basin tectonic, isostatic, stratigraphic, thermal modeling which accounts for the feedbacks between sedimentary, crustal, and mantle processes¹⁰⁶. It is designed to produce automated 2D basin models using present day stratigraphy as input, from which the structural and thermal evolution of a rifted basin is automatically reconstructed. This is achieved through the coupling of a lithosphere scale forward model with an inverse algorithm for automated parameter optimization. This kind of algorithm allows the feedbacks and interactions between shallow sedimentary and deep lithosphere/asthenosphere processes to be investigated, which are difficult to predict without iterative modeling.

Given the number of physical processes involved and the range of parameters, the evolution of sedimentary basins cannot be uniquely reconstructed. The strategy employed in TecMod is scenario analysis performed on the (user-specified) geologically most likely end member cases, which yields a possible range in target parameters (mainly temperature, heatflow and related aspects though time and space). As demonstrated by Ref.¹⁰⁶, the actual fitting process of crustal and mantle stretching factors, water depth, sediment fill, thermal conditioning and isostatic adjustment can be fully automated by casting the problem in terms of a constrained optimization problem^{107,108}. The goal function to be minimized is the misfit between observed and modeled stratigraphic data where additional constraints can be imposed (e.g. finite rifting durations, maximum stretching factors, weighting between different events). The inversion consists of an iterative search for the optimal set of crustal and mantle stretching factors and palaeo-water depth values, which minimize the chosen goal function^{109,110}. The forward run requires an initial guess for stretching factors and water depth which are successively improved by iteration with the inverse algorithm according to the scheme described by Ref.¹⁰⁶.

The 2D model is locally based on pure shear kinematics and resolves simultaneously differential thinning, flexural isostasy, heat transfer, sedimentation, erosion and compaction on a Lagrangian finite-element mesh. The general modeling strategy is that the structural and

thermal solutions are in equilibrium after every simulation time step. This is achieved by resolving simultaneously for deep lithosphere and shallow sedimentary basin processes. During every time step, the entire computational mesh is deformed with the mantle and crustal stretching factors, new sediment packages are deposited and become part of the computational domain. Flexural isostasy is ensured by further movement of the mesh in the vertical direction.

Heat transfer occurs by advection and diffusion:

$$\rho_i c_i \left(\frac{\partial T}{\partial t} + v_x \frac{\partial T}{\partial x} + v_z \frac{\partial T}{\partial z} \right) = \frac{\partial}{\partial x} \left(k_i \frac{\partial T}{\partial x} \right) + \frac{\partial}{\partial z} \left(k_i \frac{\partial T}{\partial z} \right) + A_i \quad (1)$$

where t is time and v_x and v_z are the velocities in the horizontal x - and vertical z -direction, respectively. In the employed Lagrangian reference frame, the numerical grid is deformed with the pure shear velocities so that the advection terms in equation (1) are automatically accounted for. In the case of differential thinning, re-meshing of the mantle ensures geometric consistency between the crust/sediments and the mantle portions of the modeling domain. ρ , c_i , T , k_i and A_i are density, specific heat, temperature, thermal conductivity and radiogenic heat production rate, respectively. The index i illustrates that the property is dependent on the local rock type. The thermal conductivity of the sediments is an effective material property that results from geometric averaging of the local matrix conductivity (k_m) and pore-fluid conductivity (k_w)^{111,112}:

$$k = k_m^{(1-\phi)} k_w^{(\phi)} \quad (2)$$

where ϕ is porosity. The temperature dependence of thermal conductivity is accounted for by assuming the following relations^{113,114}:

$$\begin{aligned} k_w &= c_2 T^2 + c_1 T + c_0 \\ k_m &= k_m^0 \exp\left(\frac{c_f}{T} - \frac{c_f}{T_0}\right) \end{aligned} \quad (3)$$

with c parameters defined in Table S1. Crustal and mantle densities evolve during basin formation as functions of thinning, melt production, magmatic intrusion, heat transfer and sediment and water loading as well as phase transitions^{115,116}. The resulting stresses on the lithosphere are projected onto a plate with an assumed elastic thickness and are compensated by isostasy¹¹⁷⁻¹¹⁹. The equation describing the vertical deflection of the top of the crust, w , due to flexural isostasy is:

$$\frac{\partial^2}{\partial x^2} \left(D \frac{\partial^2 w}{\partial x^2} \right) + (\rho_m - \rho_{in}) g (w + S) = q \quad (4)$$

where x , D , ρ_m , ρ_{in} , g , and S are the horizontal coordinate, the flexural rigidity, the average density of the lithospheric mantle, the average density of the basin infill (compacted sediments), the acceleration due to gravity and the kinematic deflection due to necking. The differential load, q , represents differences in load relative to the load in the equilibrium state. The flexural

rigidity, D , is a cubic function of the effective elastic thickness (T_e), which can be kept either constant or determined by the depth of a specified isotherm¹²⁰ or the thickness of the crust. The kinematic deflection S , i.e. the surface depression that results from stretching and not from isostatic compensation, is calculated using the crustal thinning factors and the depth of necking¹¹⁸. Both ends of the elastic beam are either matched to an analytical solution of a beam clamped at infinity or the modeling domain is extended laterally to minimize boundary effects. Although elastic thickness is believed to be close to zero during late rifting as the plate is thin and hot¹², flexural isostasy was nonetheless tested (e.g. Model ELT10, Table 1 in main article). The results were similar to those assuming elastic thickness is 0 km and the effect of flexure is therefore regarded as insignificant during breakup.

Sediments are incorporated in the thermal solver to resolve blanketing effects¹²¹⁻¹²⁴. Sediment deposition is controlled by the time dependent water depth which comes out of the inversion (see below). For each time step, a local sedimentation rate is iteratively computed so that at the end of a time step, i.e. after sedimentation, compaction, isostatic compensation and flexure, the basin is filled to the pre-computed water depth. The deposited sediments are compacted using empirical compaction laws, specifically those published by Refs.^{125,126} that have the form:

$$\phi_i = \phi_i^0 \exp(-c_i z) \quad (5)$$

where ϕ is porosity, c is a compaction length scale and z the burial depth.

The boundary conditions for the thermal solver are user-specified fixed temperatures at the base and top of the numerical domain and zero horizontal heat flow at the sides utilizing a finite element method-based code. This allows the use of a body-fitting Lagrangian mesh strategy and ensures that all material contrasts are accurately resolved. The model and its individual components have been tested versus a number of analytical solutions including tectonic subsidence, thermal subsidence and steady state heat flow, with radioactive heat production rates decaying exponentially with depth.

Modifications and parameterization of melting and melt migration. We have introduced new features to TecMod in order to test our main hypothesis that melt retention is an important source of buoyancy during and shortly after breakup of continental plates. This has required i) tracking of the lithosphere- asthenosphere boundary (LAB) and accounting for possible compositional changes across it, ii) resolving melt generation, iii) melt migration, and iv) melt emplacement.

Tracking of the LAB has been implemented by assigning distinct rock properties to the oceanic and continental lithospheres and asthenospheres (Table S1). The employed Lagrangian mesh allows accurate tracking of the LAB during active syn-rift and passive post-rift deformation.

Simulations of decompression melting and melt migration continue to be numerically challenging. Models explicitly designed to resolve these processes typically use published solidus functions to solve for melt generation and the equations of compaction driven fluid flow to track the produced melts during their ascent towards the surface. We have deliberately not taken this approach. Models and observations on melt generation, retention and migration are still difficult to reconcile with geophysically inferred melt fractions in various rift settings as they are often

significantly higher than predicted by models. This and the considerable technical challenges associated with implementing melt migration into basin modeling software like TecMod has led us to derive a parameterization of melting and melt migration which allows us to investigate different scenarios that are consistent with geophysically inferred melt fractions.

We have parameterized decompression mantle melting by assuming that the degree of melting scales with the heights to which the LAB rises. The onset of melting is assumed to occur at 90 km depth, consistent with published solidus functions^{65,69,94,98}. The concentration of melt is then assumed to increase linearly up to a specified depth of 20 km, where a maximum degree of melting of 16% is reached (average melt content of 8 %), which is consistent with observations at active rift zones and melt fractions calculated for adiabatic mantle upwelling⁷². Melt extraction is parameterized by assuming that the continental lithosphere is a permeability barrier for rising melts. As long as it is intact, melts are retained within the underlying asthenosphere. After breakup, loss of melt occurs exponentially at a prescribed rate of, for example, 15% every 1 My, meaning that more than half has been lost after 5 My and less than 2% melt remains in the asthenosphere after 10 My.

In addition, we have implemented the option to intrude igneous underplates. This allows us to explore scenarios in which a gabbroic sill of thickness equivalent to a specified fraction of the melt retained in the asthenosphere (e.g. 25%) is emplaced below the crust during every time step. For the Kwanza model reported in the accompanying article, this results in an underplate 5 km thick at the point of breakup tapering to 0 km towards the inboard edge of the rifted margin.

Finally, we use the models of Ref.¹¹⁶ (specifically R123 peridotite) to test the effect that mineral phase changes in the mantle (both lithosphere and asthenosphere) have on the isostatic equilibrium of the plate.

Plate reconstructions

Recent plate reconstructions for the South Atlantic (Refs.^{29,30} and Movie S4) use tectono-stratigraphic ages that fit with the comprehensive well and seismic data acquired by the oil industry in Brazil and Angola^{5,8,16}. The first oceanic crust in the central South Atlantic area is interpreted to have formed around 123 Ma (early Aptian), correlating with the breakup unconformity at the base of the pre-salt sag strata and corresponding to the age of the outer volcanic high which marks the inner edge of oceanic crust in Angola⁵. The sag strata (123-117 Ma) and the overlying salt (117-113 Ma) were deposited inboard and thinly on top of the high, which is interpreted to represent the embryonic ridge¹⁸. The favoured plate reconstructions show a relatively simple breakup (e.g. Movie S4). These reconstructions are significantly different to the models of Refs.^{24,28} where large amounts of trans-continental shearing are required in order to fit plates which remain attached for another 10 My, until the end of salt deposition. This shear movement is speculative, we believe resulting from an erroneous breakup age.

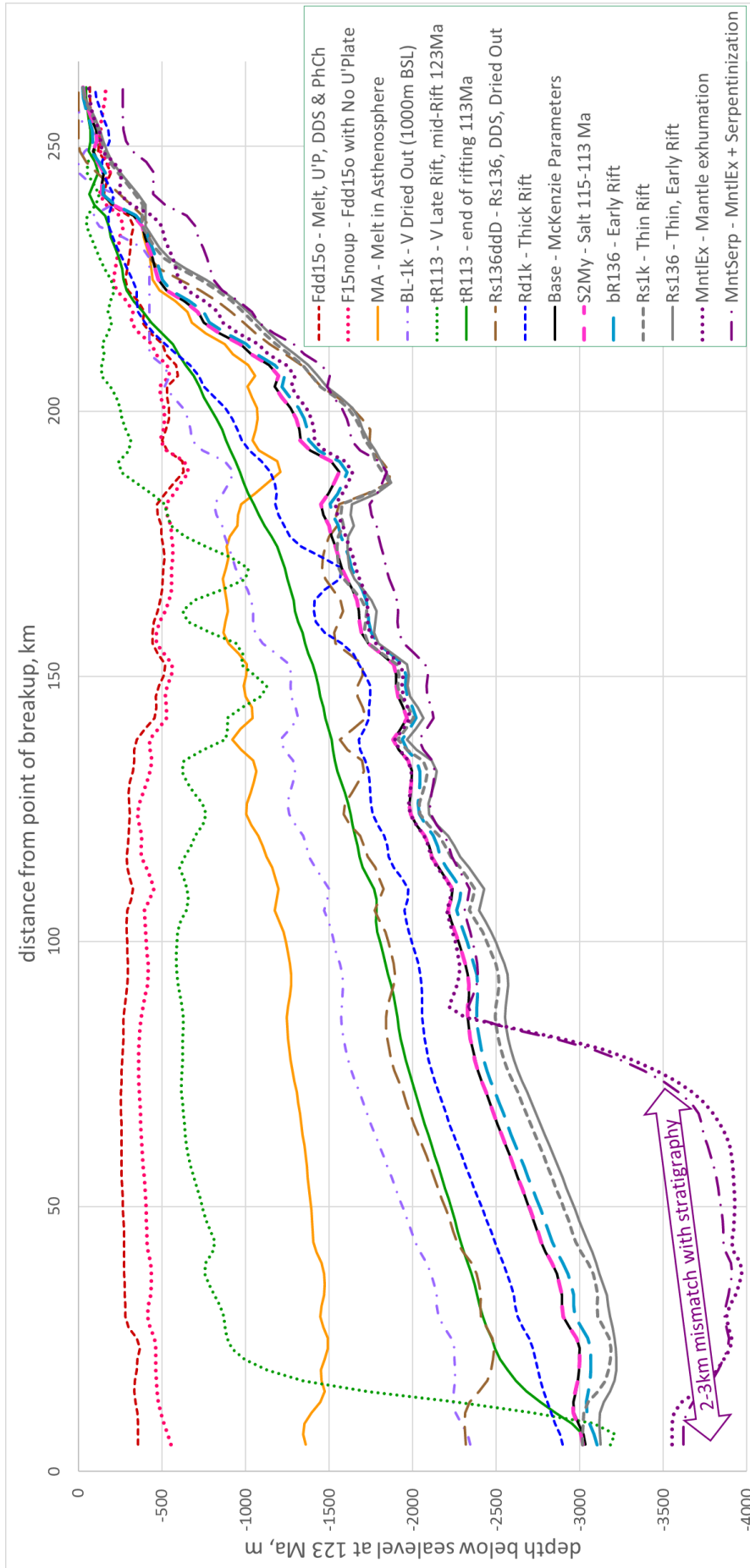


Figure S1. Depth of the syn-rift sediment surface at time of breakup, 123 Ma, for representative models in the accompanying article (Models Base, MA and Fdd15o, Fig.1a) plus additional models where the stratigraphic match is less than perfect.

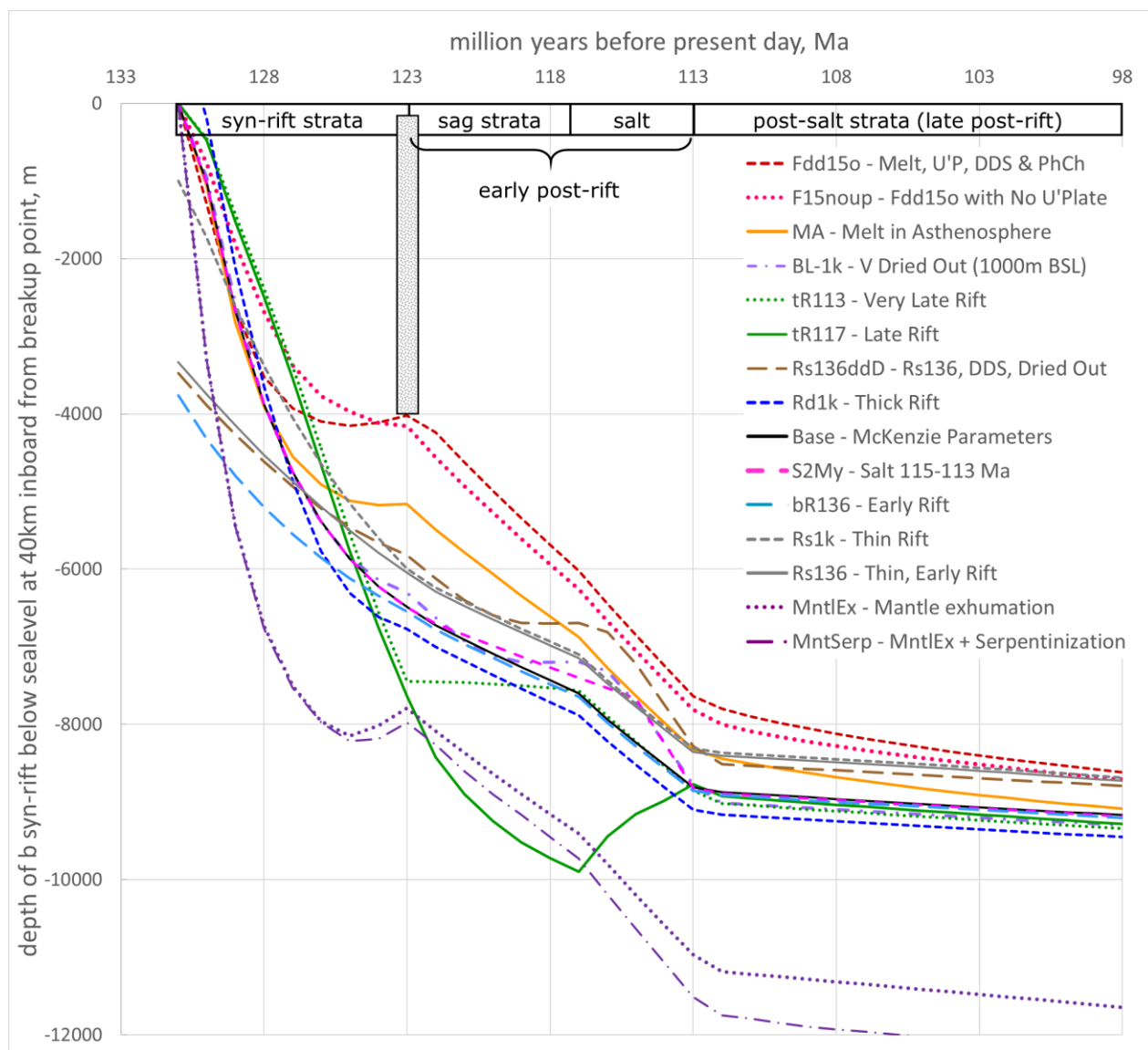


Figure S2. Graphs comparing subsidence histories of a 131 million year old surface (base of rift) for the models shown in Figure S1 at a position 40 km inboard of the breakup point. The grey bar indicates the thickness of syn-rift sediment at the time of breakup. Differences between the grey bar and a specific subsidence line at 123 Ma represents the water depth (unfilled accommodation space). Palaeobathymetric evidence of contemporaneous erosion and subsequent shallow water sag and salt deposits in the central South Atlantic region is contradicted where the water depth is significant (>500 m).

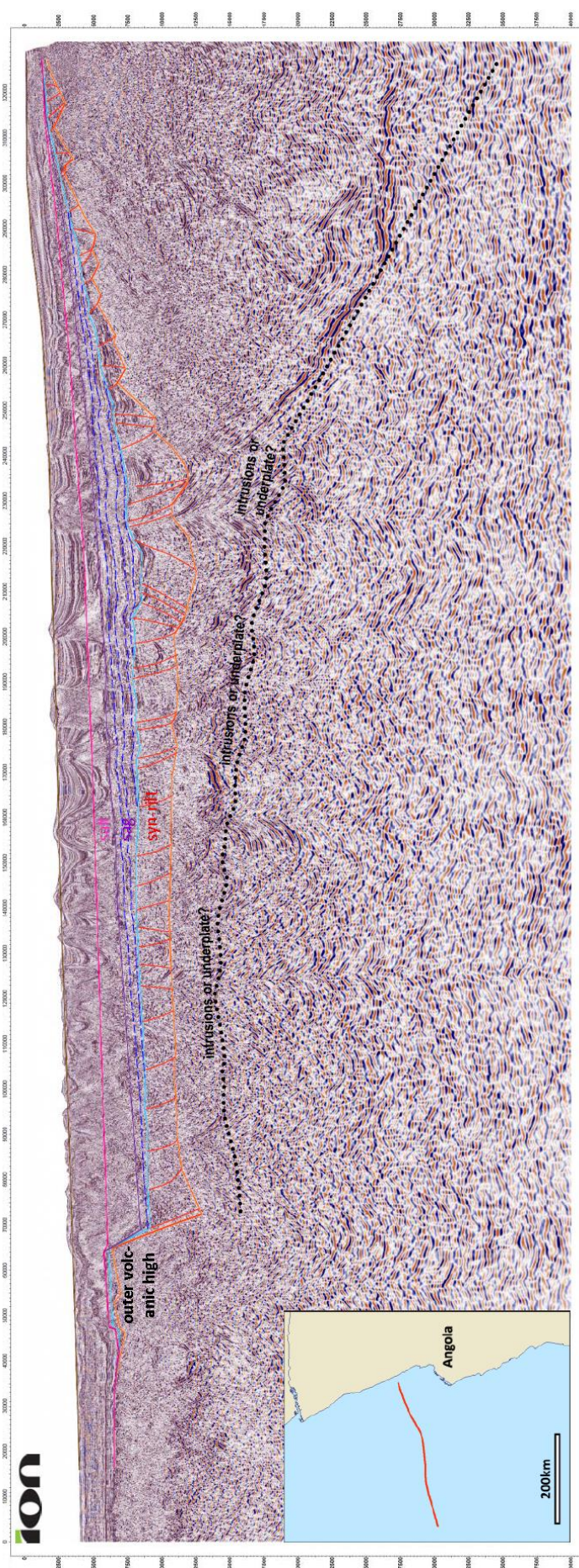


Figure S3. 2D pre-stack depth migrated seismic section from northern Kwanza Basin, offshore Angola, including the interpretation used as input for modeling of a non-volcanic rifted margin reported both here and in the accompanying article. Data is courtesy of ION GeoVentures (CongoSpan Seismic Line 2400). The field of view is 330 km (horizontal) by 40 km (vertical): the scales shown are in metres. West is on the left-hand (ocean) side of the section, east on the right (continent). Solid lines are the interpreted horizons used in TecMod, excluding indicative faults shown in red. Dashed dark blue lines indicate that flat layers within the sag progressively step (lap) onto the erosional surface at the top of the syn-rift interval (light blue line). The black dotted line indicates the base of crust which is an output from one of our final models (Fdd15o), its position correlating with the strong reflections generally interpreted to mark the Moho. The data was originally interpreted in Petrel 2016 and then exported as an image to Adobe Illustrator CC 2015 where the horizons were redrawn. The interpretation was then imported into TecMod as an SVG file to carry out the modeling. Brown horizon – seabed, 0 Ma; pink horizon – top salt (pre-salt diapir geometry), 113 Ma; purple horizon – base salt/top sag, 117 Ma; blue horizon – base sag/top syn-rift, 123 Ma; orange horizon – base syn-rift/top basement, 131 Ma. Oceanic crust starts at the outer volcanic high which we interpret to be a proto-oceanic ridge formed at the time of breakup, 123 Ma. The inset map shows the location of the section (red line). The same section is shown in Figs. 5 and 7 in Ref.¹⁰ as well as others noted in the same reference.

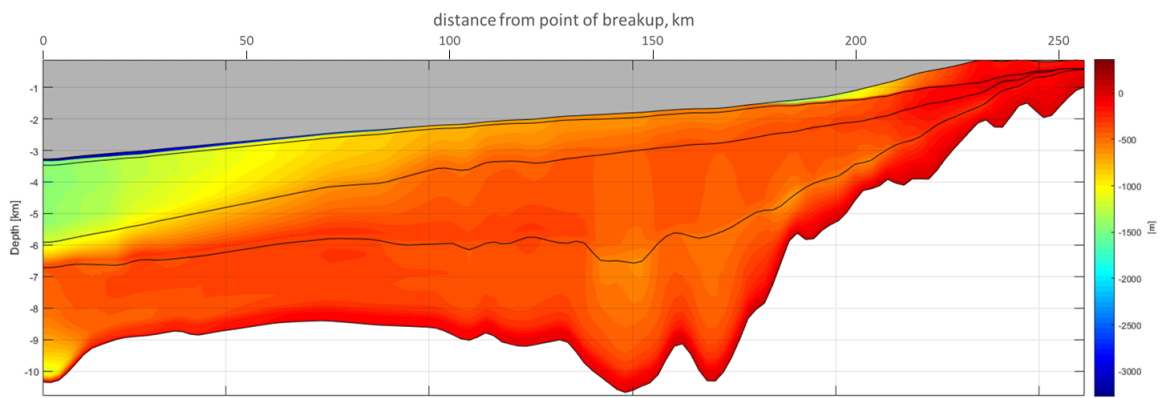


Fig. S4a

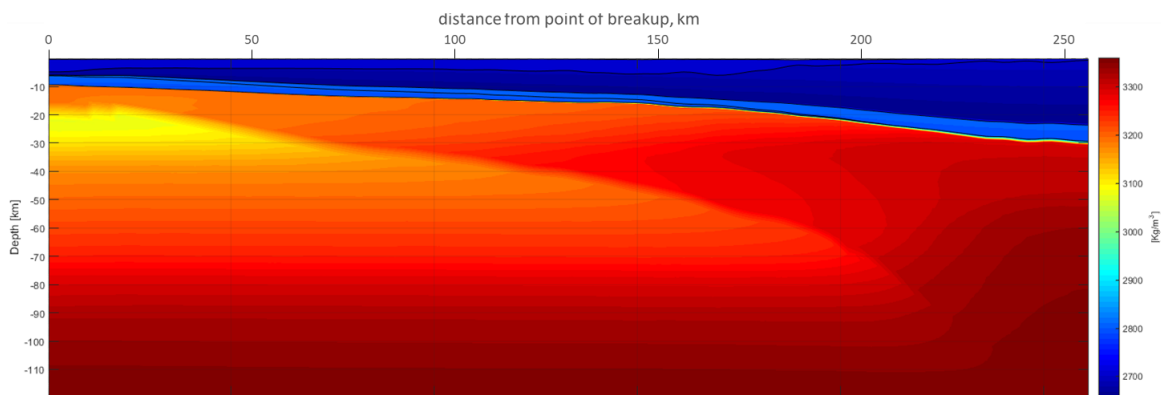


Fig. S4b

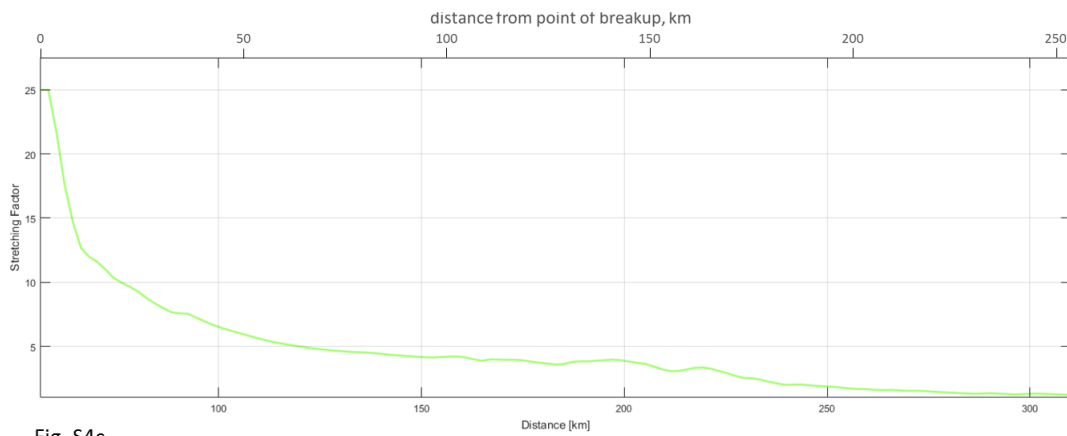


Fig. S4c

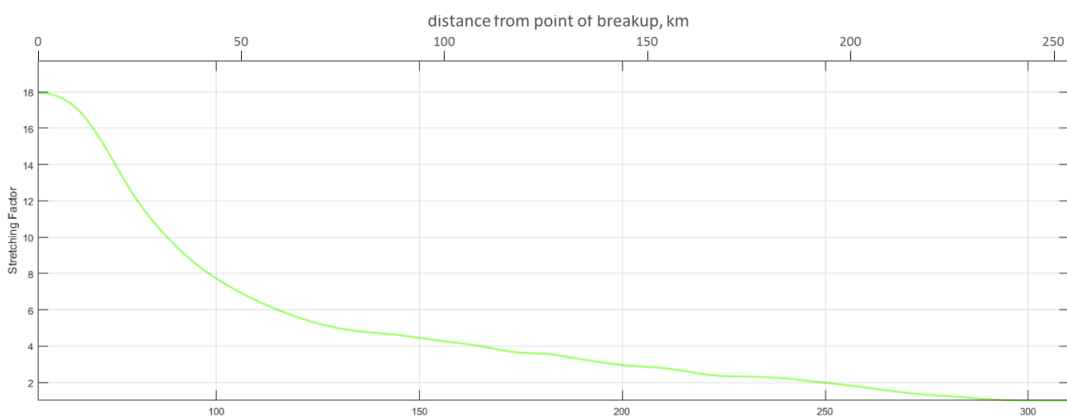


Fig. S4d

Figure S4. Cross-sections from TecMod model Fdd15o, supplemental to those shown in Fig.1 in the main article. **(a)** Model palaeobathymetry (depositional depth of sediments) shown in color. The scale is shown on the right, with water depth in metres. **(b)** Density of crust (blue colors) and density of asthenosphere (red-yellow colors) at the time of breakup. The scale is shown on the right, with densities in kg.m^{-3} . **(c)** Stretching factor for the crust (not including thickening by underplates). **(d)** Stretching factor for the mantle lithosphere.

<u>Parameter</u>	<u>Description</u>	<u>Value</u>	<u>Units</u>
	initial upper crustal thickness	28.0	km
	initial lower crustal thickness	7.0	km
	initial lithosphere thickness	120	km
	isostatic compensation depth	120	km
	initial height of land surface	500	m (above sea-level)
	basal lithosphere temperature	1603	°K
	seafloor temperature	288	°K
ρ	density		kg/m ³
c_i	specific heat capacity		J/kg/°K
T	temperature		°K
t	time		s
x	horizontal coordinate		m
z	vertical coordinate		m
v_x, v_z	horizontal and vertical velocity		m/s
A	radiogenic heat production in crust	2	$\mu\text{W}/\text{m}^3$
k	thermal conductivity		W/m/°K
α	thermal expansion		1/°K
ϕ	porosity		decimal value
c	inverse compaction		km ⁻¹
c_0	parameter in pore fluid thermal conductivity	0.565	W/m/°K
c_1	parameter in pore fluid thermal conductivity	1.88E^{-03}	W/m/°K ²
c_2	parameter in pore fluid thermal conductivity	-7.23E^{-06}	W/m/°K ³
c_f	parameter in grain thermal conductivity 250		°K
D	flexural rigidity		Pa.m ³
g	gravitational acceleration		m/s ²
w	vertical deflection of the crust		m
S	kinematic deflection		m
q	differential load		kg/m ³
z	burial depth		m
	density of melt	2650	kg/m ³

Table S1. Model parameters.

<u>Stratigraphic Name</u>	<u>Age Ma</u>	<u>Lithology</u>	<u>ϕ_i</u>	<u>c km⁻¹</u>	<u>α 1/°K</u>	<u>ρ kg/m³</u>	<u>k W/m/°K</u>	<u>c_i J/Kg/°K</u>	<u>A W/m³</u>
post-salt	0	shale	0.62	0.50	0	2720	1.5	1000	1 x 10 ⁻⁶
salt	113	halite	0.05	0	4.4 x 10 ⁻⁵	2300	1.2	1000	1 x 10 ⁻⁶
sag	117	90% sh, 10% lst	0.60	0.49	0	2710	1.5	1000	1 x 10 ⁻⁶
syn-rift	123	50% ss, 50% clst	0.56	0.39	0	2705	2.6	1000	1 x 10 ⁻⁶
upper crust	131	granite	0	0	2.4 x 10 ⁻⁵	2700	2.5	1116	3 x 10 ⁻⁶
lower crust		mafic	0	0	2.4 x 10 ⁻⁵	2850	2.5	1116	2 x 10 ⁻⁶
mantle lithosp		peridotite	0	0	3.2 x 10 ⁻⁵	3300	3.4	1288	0
asthenosphere		fertile peridotite	0-0.16	0	3.2 x 10 ⁻⁵	3300	3.4	1288	0

Table S2. Physical properties of rock types used in the models. Sh = shale; lst = limestone; ss = sandstone; clst = claystone.

Movie S1. 2D TecMod reconstruction for model Fdd15o through the upper part of the lithosphere showing the evolution of crust and sediment infill from 131 Ma to present day. Open with an internet browser. Colors indicate rock type. Pale blue = water.

Movie S2. 2D TecMod reconstruction for model Fdd15o through the lithosphere and upper part of the asthenosphere showing the melt fraction, indicated by colors, from 131 Ma to 91 Ma. Open with an internet browser.

Movie S3. 2D TecMod reconstruction for model Fdd15o through the lithosphere and upper part of the asthenosphere showing the density, indicated by colors, from 131 Ma to present day. Open with an internet browser.

Movie S4. Plate reconstruction for South Atlantic region based on work reported in Ref.⁵. Open with Microsoft PowerPoint, use Slide Show and the animation starts automatically on slide 2. The plate model was built using PaleoGIS software with the main assumptions shown on slides 15 and 16.

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