1 Shallow seismic investigations of the accretionary complex offshore Central Chile

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10 Abstract

11 Thrust ridges are accretionary structures often associated with local uplift along splay faults and cold seep 12 activity. We study the influence of a NS-trending thrust ridge system on the transition between the accretionary 13 prism and the continental framework (shelf break) offshore the Maule Region (central Chile at 35°-36°S) by 14 examining its 2-D and 3-D seismic velocity structure. The experiment comprises five densely spaced seismic 15 refraction lines running subparallel to the trench and recorded at nine OBH/S (ocean bottom hydrophone/seismometers) deployed along the central line. Results show a narrow margin-parallel volume 16 17 (approximately 6x50x5 km³) whose velocity distribution is consistent with sedimentary rocks. The shallow 18 sedimentary unit is characterized by the presence of very low velocity hydrate-bearing sediments (<1.7 km/s), 19 which are interpreted as highly porous sedimentary rocks (> 50% porosity) within the Gas Hydrate Stability Zone 20 (GHSZ) suggesting low hydrate content. These zones spatially correlate with fluid activity in the vicinity of the 21 NS trending thrust ridges based on local high heat flow values (>40 mWm⁻²) and seepage mapping. On the other

- 22 hand, the splay faults that crop out on the flanks of the thrust ridge structures might be responsible for tectonically
- 23 induced vertical fluid migration.

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25 Key Words: Accretionary Prism; fluid migration; Chile; splay fault; sediment

26 1. Introduction

27 Accretionary convergent margins are typically characterized by a slow convergence rate (<7.6 cm/year) 28 and trench sediment thickness greater than 1 km (Clift and Vannucchi, 2004). This is the case of the Nankai (Moore 29 et al., 1990), Cascadia (Hyndman et al., 2011), and south-central (SC) Chile (Contreras-Reyes et al, 2010) margins. 30 Trench sediments can be accreted frontally (forming a large accretionary prism with low slope angle), or basally 31 (causing oversteepening of the continental slope) (van Gool and Cawood, 1994; Contardo et al., 2008; Contreras-32 Reves et al., 2016). In particular, the SC Chilean continental margin currently falls within the classification of an 33 accretionary margin: a slow convergence rate (6.6 cm/year; Angermann et al., 1999), and trench sediment 34 thickness of ~2 km (Díaz-Naveas, 1999; Grevemever et al., 2003; Maksymowicz et al., 2017; Contreras-Reves et 35 al., 2017; Tréhu et al., 2019). However, the SC Chilean margin underwent a period of subduction erosion during 36 the Miocene, when the convergence rate was > 8 cm/year and the trench sediment thickness was < 1 km (Kukowski 37 and Oncken, 2006). A rapid increase of sediment supply to the trench in the Pliocene following a period of 38 glaciation and a steady decrease of the convergence rate between the Nazca and the South American plates shifted 39 the SC Chilean margin from erosive to accretionary (Melnick et al., 2006).

40 Seismic studies have reported the presence of an accretionary prism 30-50 km wide with P-wave velocities (V_p) of 2.5-5.0 km/s interpreted as unconsolidated (frontal prism) and semi-consolidated accreted (middle prism) 41 42 sediments (Fig. 1; Moscoso et al., 2011; Contreras-Reyes et al., 2017). The landward edge of the interpreted 43 accretionary prism is defined by an abrupt increase in seismic velocities and the location of the shelf break, defined 44 as the transition from the continental slope to the shelf (Fig. 1; Contreras-Reves et al., 2017; Tréhu et al., 2019a). 45 The transition between the accretionary prism (frontal and middle prisms) and the continental framework (internal 46 prism) is usually referred as the backstop and has been proposed to mark the up-dip limit of the seismogenic zone, 47 inhibiting earthquake rupture propagation towards the trench (Contreras-Reyes et al., 2010; Moscoso et al., 2011). 48 However, evidence for seafloor uplift along a profile crossing the patch of greatest slip during the 2010 Maule 49 earthquake suggests at least localized up-dip activation of the interplate boundary beneath the accretionary prism during mega-thrust earthquakes, with important implications for the processes of tsunami generation and
rheological behaviour of the continental wedge (Maksymowicz et al., 2017).

52 Generally, accretionary margins consist of a series of mostly trench-parallel accretionary thrust ridges that 53 contain largely compressed, folded, and faulted turbidites and trench-fill deposits (Suess et al., 2018). These 54 morphological features, characterized by relative seafloor uplift of up to hundreds of meters, have been observed 55 in accretionary prisms worldwide such as in the Barbados Ridge accretionary complex (e.g. Brown et al., 1987), 56 the Nankai accretionary complex (e.g. Park et al., 2002; Schumann et al., 2014), the southwestern Taiwan margin 57 (e.g. Klaucke et al., 2016) and Cascadia (e.g. Hyndman et al., 1994; Tréhu et al., 1995, 1999; Fisher et al., 1999; Adam et al. 2004). The folding and thrusting around the accretionary ridges are usually associated with splay 58 59 faulting across the trenchwardmost edge of the accretionary prism. Bottom simulating reflectors (BSR) are often 60 present within these accretionary ridges and are associated with accumulation of free gas underneath possible gas 61 hydrate formation, which might be interrupted by cold seep activity (seepage) (Brown et al., 1987; Tréhu et al., 62 2004a, 2004b; Klaucke et al., 2016). Mud diapirism and volcanism can occur as well in these compressional 63 scenarios providing an important dewatering mechanism of the accretionary wedge (Brown et al., 1988; Kopf, 64 2002). Thrust faults may also act as dewatering pathways (Cloos, 1984; Klaucke et al., 2016).

65 Accretionary ridges have been found in the accretionary prism along the Chilean margin (e.g. across the 66 accretionary prism off Golfo de Arauco and Penas; Contreras-Reyes et al., 2008 and Maksymowicz et al., 2012; 67 respectively). Near the shelf break around the Maule Region, Geersen et al. (2011) identified accretionary thrust 68 ridges that probably correspond to the superficial expression of active splay fault systems. In this study, we image 69 the subsurface beneath a previously mapped (Geersen et al., 2011) N-S trending thrust ridge (Fig 3) that has been 70 subject of investigation of small-scale structures like slumps, chemoherms, mud volcanoes, gas hydrate deposits, 71 fluid migration path and seep sites (Grevemeyer et al., 2003; 2006; Flueh and Bialas, 2008; Klaucke et al., 2012; 72 Villar-Muñoz et al., 2014). We present a shallow margin-parallel V_p model of the landward edge of the 73 accretionary prism near the shelf break offshore SC Chile (35°-36°S) based on controlled source seismic data. To 74 evaluate the robustness of our results, the data were analysed using two different travel-time modelling techniques. 75 We image the 3D P-wave structure of the upper ~4 km in the transition zone between the accretionary prism and

76 the continental framework and discuss the implications of the model for local tectonics. Our results provide new



77 insights into the interplay between the dynamics of accretionary structures and local distribution of fluid activity.

Fig. 1. Interpretative summary sketch of the southern-central Chilean margin for the Maule segment based on seismic studies (Moscoso et al., 2011; Contreras-Reyes et al., 2017). The accretionary prism is composed by the outer/frontal prism (poorly consolidated sediment) and middle prism (more compacted and lithified sediment). The continental framework rock (or inner prism) corresponds to the paleo-accretionary prism or continental basement (Contreras-Reyes et al., 2010). The solid dark green line indicates the accretionary Prism/Continental Framework Rock contact or backstop. Our study is located near the shelf break at the transition between the middle prism and the inner prism.

85 2. Tectonic Setting

86 The SC Chilean margin is controlled by the subduction of the oceanic Nazca Plate under the continental 87 South American Plate at a current rate of 6.6 cm/year and a convergence azimuth of about 78° (Fig. 2; Angermann 88 et al., 1999). The trench is filled with sediments 1-2.5 km thick (Völker et al., 2013) and is incised by a sinuous 89 axial channel. The current accretionary regime (Bangs and Cande, 1997) is evident by the morphology of the 90 forearc westernmost part which is characterized by a 30-50 km wide accretionary prism that can be subdivided 91 into two main segments: (1) an outer accretionary wedge (5-10 km wide) characterized by low seismic velocities 92 of <3.0 km/s, and interpreted as a frontal prism of poorly compacted and highly deformed sediments, and (2) the 93 middle wedge (~50 km wide and V_p of ~4 km/s) interpreted as a middle prism composed by older, compacted and 94 lithified sediment, overlain by an apron of low velocity slope sediments (Fig.1; Moscoso et al., 2011; Contreras-95 Reves et al., 2017).

96 Our study area is located at the transition between the middle prism and the continental framework rock 97 (Figs. 1-3). Seismically, this boundary has been characterized at depth by an abrupt lateral change in velocity from 98 velocities characteristic of accretionary prism sediments (<4 km/s) to velocities of >5 km/s representing the 99 continental framework rocks. The continental framework in this region has been interpreted to be a late 100 metamorphosed Paleozoic subduction complex (Willner, 2005). Splay faults can be generated and intersect the 101 seafloor forming thrust ridges in this transition zone in response to compression and shortening of the forearc 102 (Moore et al., 2007). Formation of intra-slope basins with hemipelagic sediments is favoured by the existence of 103 these ridges (Suess et al., 2018). In our study area, a thrust ridge system and associated intra-slope basins have 104 been mapped (dashed pink line in Figs. 2, 3) and interpreted by Geersen et al. (2011) as the seafloor expression of 105 the splay fault accommodation of subduction motion. Geersen et al., (2011) propose this thrust ridge as the 106 boundary between the highly deformed accretionary prism and the continental shelf sedimentary cover overlying 107 the continental framework.

Much of the SC Chilean margin is characterized by basal accretion and sediment underthrusting, e.g. at the NW of our seismic survey as evidenced by the seafloor morphology of the accretionary prism near the trench ("Reloca slide" in Fig. 3; Contreras-Reyes et al., 2016) and subduction channels of considerable thicknesses (Olsen et al., 2020). However, both frontal accretion and sediment underthrusting have also been observed (ContrerasReyes et al., 2010; 2017; Tréhu et al., 2019a). Another bathymetric feature observed in this area corresponds to
the Itata Canyon (Figs. 2, 3A), which incises the continental shelf; however, this system is not present at the lower
continental slope (Geersen et al, 2011) and is thought to be inactive at present (Klaucke et al., 2012). Tréhu et al.
(2019a) discuss tectonic implications of similar canyons (the Mataquito and Huenchullami canyons) that do not
reach the trench between 34° and 35°S and attribute this to large scale disruption of the forearc due to subduction
of topography.



Fig. 2. South-central Chile bathymetry from the Global Multi-Resolution Topography (Ryan et al., 2009), version 3.6.6. Main
 oceanic bathymetric features are depicted. Convergence between the oceanic Nazca and the continental South American Plates

- 121 is indicated by the yellow arrow. Red contoured square encloses the location of the study area shown with more detail in Fig.
- 3. See the legend for further information.

123 Several authors have proposed the existence of fluid activity in accretionary ridges in SC Chile. Klaucke 124 et al., (2012), using sidescan sonar mapping and seafloor observations, reported the occurrence of both fossil and 125 likely active cold seeps. As suggested, chemoherms (carbonate structures) and buried seeps encountered in our 126 study area were partially formed or have been overprinted by fluid venting and biogenic formation probably 127 generated within the slope sediments (Fig. 3). However, Klaucke et al. (2012) propose little current activity, 128 although either much higher fluid fluxes in the past or fluid flows over prolonged periods may explain the 129 observations. On the other hand, bottom simulating reflectors (BSR) are often present in accretionary margins and 130 are commonly associated with the occurrence of gas hydrates on the continental slope. Heat flow anomalies 131 derived by BSR depth calculation of geothermal gradients have been obtained in the vicinity of our study area 132 (Grevemeyer et al., 2003; Villar-Muñoz et al., 2013) with values >35 mWm⁻², indicating advective fluid migration 133 along stratigraphic boundaries or fault zones (Fig. 3). Grevemeyer et al. (2006) calculated heat flow anomalies 134 from ODP Leg 202 drillcore data along a transect that intersects our profiles with values up to ~120 mWm⁻² (Fig. 135 3).

136 Seismic reactivation of splay faults occurred 12 days after the 2010 Maule megathrust earthquake $M_w 8.8$ 137 in the Pichilemu area (34.5° S; Farías et al., 2011; Ruiz et al., 2014). Using a local marine network, some seismic 138 activity was detected mainly near the boundary between the active accretionary prism and continental basement 139 as a post-seismic response of the outer accretionary wedge updip from the patch of greatest slip during the 2010 140 M_w 8.8 Maule earthquake (Tréhu et al., 2019b). In contrast, little post-seismic activity was observed in the 141 accretionary prism and along the splay faults around our study area (See Fig. 1 in the Supplementary material, 142 section A1; Moscoso et al., 2010; Lange et al., 2012; Contreras-Reves et al., 2017). Lieser et al. (2014) proposed 143 the existence of a stable accretionary prism to explain the progressive decrease in post-seismic splay fault response 144 south of $\sim 36^{\circ}$ S.

145 **3. Seismic Modelling**

146 3.1. Wide-Angle Seismic Data

147 Seismic data were acquired by the German IFM-GEOMAR Institute for Marine Sciences (now GEOMAR 148 Helmholtz Centre for Ocean Research Kiel) during cruise JC23 on the RRS JAMES COOK in 2008. The 149 experiment comprised five parallel wide-angle seismic profiles spaced ~1.5 km apart with strike of N31.5°E and 150 maximum length of 50 km (Fig. 3). Spatially coincident deep towed sidescan data were also acquired. The source 151 was 3 GI airguns each equipped with 250 cu-inch generator and 105 cu-inch injectors were shot at 13 sec intervals 152 providing an average shot spacing of ~20 m (Flueh and Bialas, 2008). Along the central line, 11 OBS/H (Ocean 153 Bottom Seismometers/Hydrophones) were deployed on the seafloor at a depth of ~1500 m. However, due to low 154 signal to noise ratios on the geophones, we only used data from the 3 OBHs and from the hydrophone component 155 of 6 OBSs for further processing and modelling (Fig. 3). A 4-channel surface streamer (300 m long) also recorded 156 the shots.



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159 Fig. 3. Study area, experiment geometry, main bathymetric and tectonic features, and seepage activity. The high-resolution 160 bathymetric image of the seafloor off Maule region was collected during the cruise JC23 on the RRS JAMES COOK 161 (Flueh and Bialas, 2008). Red and orange circles indicate the 9 used stations and thin black lines correspond to the wide-angle 162 seismic refraction profiles (L10, L11, L12, L13, L14) processed in this study whose gunshot locations are numbered every 163 250 shots (black dots along the profiles). Solid thick lines depict locations of P09 (Contreras-Reyes et al., 2016) and ENAP-164 1 profiles (Geersen et al., 2011; see Fig. 2 for the whole extent). The N-S/NE-SW trending Thrust Ridge (dashed pink line) 165 is clearly visible by its bathymetric signature whose location coincides with thrust splay faults visible along ENAP-1 profile 166 (Geersen et al., 2011). Fluid expulsion indicators and their locations are also shown.

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Seismic data examples recorded at four instruments from shots along profile L10 are shown in Fig. 4. A strong hyperbolic signal centred on the instrument corresponds to the water wave arrival. At greater offsets, first arrivals were identified and classified into two groups based on their distinct slope: low velocity sedimentary phases (P_s) and deeper sedimentary phases (P_g). The character of the data shows a strong triplication suggesting
an abrupt increase in velocity at a boundary between the layer producing phases P_s and the one producing P_g.

A total of ~13,000 first arrivals were picked. Direct water wave arrivals were not picked because acquired high-resolution bathymetry constrain the seafloor depth and pre-processing aboard the cruise was already done in order to relocate the instrument positions. Pick uncertainties were estimated to be 60 ms. Experiment geometry and poor signal propagation limited the picking of arrivals with offsets longer than 15 km, limiting the maximum depth of penetration to ~4 km.



- 179 Fig. 4. Seismic record examples along profile L10 recorded at (A) OBH 1002, (B) OBH 1004, (C) OBS 1005 (respective
- 180 hydrophone channel) and (D) OBS 1011 (respective hydrophone channel). For each seismic record, the upper panel shows
- 181 the corresponding travel-time section using a reduction velocity of 6 km/s and the detected first-arriving P waves refractions
- 182 are indicated: P_w (water wave arrival), P_s (low velocity sedimentary phase) and P_g (deeper sedimentary phases). Corresponding
- 183 picked phases (P_w excluded) and respective uncertainties are shown in the middle panel: yellow dots denote P_s phases and red
- 184 dots denote P_g phases. Lower panel shows the location of the instruments along profile L10. For seismic records examples
- along profiles L11, L12, L13 and L14, see Fig.2 in the Supplementary Material, Section A2.

186 *3.2. Three-Dimensional Travel-time Tomography*

187 Travel-time inversion and 3-D ray tracing in a heterogeneous media was computed using the tomographic 188 software package TOMOLAB/STINGRAY (Toomey et al., 1994). A 31.5° rotated Cartesian coordinate system 189 was used for defining a slowness model, which was parameterized as a $10 \times 54 \times 8 \text{ km}^3$ gridded volume with a 190 cell size of 100 m in all directions. A minimum-structure initial model was constructed through 3-D extrapolation 191 by hanging a 1-D velocity depth profile from the seafloor (Fig. 5). This 1-D model was derived by 1-D forward 192 modelling of travel times along the central profile L10 and is characterized by V_p values of 1.6 km/s at the seafloor.



Fig. 5. Experiment geometry and 2D example slices of the 3D initial velocity model used for travel-time inversion using
STINGRAY/TOMOLAB software package (Toomey et al., 1994). (A) Experiment setting using Cartesian coordinate

projection with 31.5° rotation angle with respect to north. Small black dots correspond to airgun shots along profiles L10,
L11, L12, L13 and L14. Red dots represent the 9 seismic stations used in this study. Yellow lines indicate the location of the
slices depicted in (B) and (C). (B) Initial velocity model along Y axis at X=0 km. (C) Initial velocity model along X axis at
Y=0 km.

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This inversion approach allows the definition of a perturbational model with its associated uncertainty and the subjective choice of several inversion parameters (Toomey et al., 1994). The ill-conditioned inversion problem requires a regularization which minimizes an objective function for the 3D iterative process that penalizes the roughness and size of the slowness perturbations with respect to the 3D starting model. Inversion constraints (penalty function, horizontal and vertical smoothing) are applied to the perturbational model and then normalized to remove bias (Toomey et al., 1994).

In this case, we defined a uniform node spacing of 200 m and maximum perturbation at each node of 0.05 (fractional uncertainty) for the perturbational model. We assumed an isotropic medium due to limited azimuthal ray coverage and adopted a "jumping strategy" for the 3-D tomographic inversion of first arrivals. Thus, the model constraints are operative on the total perturbation expressed with respect to the starting model (Toomey et al., 1994). implying a rather conservative approach. Several tests were computed (~50 tests) to obtain an appropriate data misfit.

The final 3-D model was obtained using a penalty function λ_p of 0.01 and vertical and horizontal smoothing factors (λ_z , λ_{xy} , respectively) of 200 and 400, respectively, resulting in an RMS data misfit of 38 ms after five iterations. Each iteration comprised a total of 3,000 LSQR iterations. A shallow ray coverage, which is a function of the velocity structure as well as of the acquisition geometry, results in shallow imaging (< 6 km); however, several model features are identified and give insights into the relationship between the velocity structure and local tectonics.

220 *3.3. Results*

221 3.3.1. Three-Dimensional Results

222 Figs. 6 and 7 depict vertical model slices along the Y and X axes, respectively, and Fig. 8 shows 223 subhorizontal slices parallel to the seafloor (along Z axis). The final velocity model is masked using a spatially 224 averaged DWS (Derivative Weight Sum) calculated as the weighted sum of all ray path lengths influenced by any 225 model parameter (Toomey et al, 1994). Overall, observed P-wave velocities in the range of 1.7-4 km/s are 226 consistent with the sedimentary nature of the middle accretionary prism (Moscoso et al., 2011; Contreras-Reyes 227 et al., 2017). The obtained velocity gradient is smoother than expected for the abrupt boundary between the 228 sediments and underlying material with velocity ~4 km/s suggested by the strong triplication of the seismic 229 records. Near the seafloor, lower velocities (<1.7 km/s, shown in cyan colours in Figs. 6, 7 and 8) are observed 230 mainly in the south-west area, which corresponds to the thrust ridge location identified by Geersen et al. (2011). 231 Velocities < 1.5 km/s (even lower than 1.4 km/s, depicted as white spots within ray-covered areas in Figs. 6, 7 and 232 8) have been observed by several authors in similar shallow hydrate-bearing sedimentary environments (Arsenault 233 et al., 2001; Hornbach et al., 2003; Schumann et al., 2014) and interpreted to indicate the presence of free gas in 234 the sediment pore space. However, several tests are performed to assess if the 3D model is capable of resolving 235 these features with a special focus on the central line (see Figs. 9 and 10 for uncertainty estimates and resolution 236 tests). In contrast, both the near-seafloor velocity (> 1.9 km/s) and the velocity gradient are larger to the northeast 237 (with a value of $\sim 1.0 \text{ s}^{-1}$), which may reflect a less perturbed sedimentary sequence over a more compact and 238 lithified accretionary wedge. The highest V_p values (~5 km/s) are present to the northeast at Y=5-10 km and a 239 depth of 4 km below the sea surface, although it is poorly resolved (Fig. 6).

The first-order correlation between lower shallow sedimentary P-wave velocities and the thrust ridge
location extends across the entire imaged volume (Figs. 6, 7 and 8). Besides, stronger velocity gradients below
Z=1 km and Y > 0 km (seen in Figs. 6, 7D, 7E, 8C and 8D) are also correlated to the associated splay fault system
imaged by Geersen et al., (2011) along the ENAP-1 profile (Fig.3).



Fig. 6. Vertical slices through the final 3D velocity model along Y axis. (A) Experiment setting using Cartesian coordinate projection with 31.5° rotation angle with respect to north. Small black dots correspond to airgun shots along profiles L10, L11, L12, L13 and L14. Red dots represent 9 seismic stations used in this study. Velocity model slices at (B) X= -1.5 km, (C) X = 0 km. (D) X = 1.5 km and (E) X = 2.3 km, are shown and their map location is indicated by yellow straight lines in (A). Dashed lines indicate iso-surfaces along Z axis which is defined relative to the seafloor.



Fig. 7. Vertical slices through the final 3D velocity model along X axis. (a) Experiment setting using Cartesian coordinate projection with 31.5° rotation angle with respect to north. Small black dots correspond to airgun shots along profiles L10, L11, L12, L13 and L14. Red dots represent 9 seismic stations used in this study. 2D velocity model slices at (B) Y= -10 km, (C) Y = -5 km. (D) Y = 5 km and (E) Y = 10 km, are shown and their map location is indicated by yellow straight lines in (A). Dashed lines indicate iso-surfaces along Z axis which is defined relative to the seafloor.



Fig. 8. Seafloor-parallel slices through the final 3D velocity model at (A) Z=0 km, (B) Z= 1 km, (C) Z= 2 km, (D) Z = 3 km

and (E) Z = 4 km, where Z is positive downwards, starting from the seafloor. These iso-surfaces are also shown in the vertical slices depicted in Figs 6 and 7. The coordinate system coincides with that of Figs. 5-7.

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261 To assess the accuracy and resolution of the 3D model, we performed a Monte Carlo uncertainty analysis 262 and synthetic recovery tests. The Monte Carlo analysis was carried out by generating several perturbed initial 263 velocity models (±5% amplitude applied to the starting velocity model) and corresponding travel time sets with 264 added random noise (±100 ms). Running corresponding inversions with the same parameters of our final model 265 allowed us to compute the average and standard deviation of the output ensemble. Results are shown in Fig. 9. 266 Areas with dense ray coverage (e.g., right below the instruments and in the middle portion of the south-western 267 part of the model), show lower uncertainty values (<0.01 km/s) whereas the north-eastern part of the model shows 268 comparatively higher uncertainty values (>0.01 km/s) (Fig. 9A, 9C). Standard deviation values show a general 269 increase with depth and towards the edges of the model, reaching maximum values around the high velocity zone at around 35 km distance (~0.2 km/s). Over most of the very low velocity zones with V_p < 1.5 km/s (white spots 270 271 in the average velocity model shown in Fig. 9B), the uncertainty values are higher than 0.02 km/s, implying a 272 possible underestimation of those velocities within that range.



Fig. 9. Monte Carlo tests for the 3D velocity model. (A) Standard deviation model along the seafloor-parallel slice at Z=0
km, i.e., at the seafloor. (B) Average velocity model at the seafloor. (C) Standard deviation on a vertical slice at X=0 km,
roughly coincident with central seismic profile L10. Distance in km is calculated from the negative end of the Y-axis in (A)
and (B), that is, an offset of 27 km is considered. (D) Average velocity model from the vertical slice shown in (C).

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Pseudo checkerboard tests were done by superimposing elliptical Gaussian anomalies of a ±5% maximum amplitude on the final velocity model across the vertical slice depicted in Fig. 6C, which is coincident with the central line L10 (see Fig. 3 for profile location). These anomalies were then prolonged along the X-axis crossing the entire 3D volume as horizontal elliptical cylinders. The distribution and sizes of the anomalies across the central vertical slice are shown in Fig. 10A. The recovery results depicted in Fig. 10B show that, as expected, the resolution is best near the central part of the model and decreases towards the edges of the model. Although most of the anomalies are recovered, significant vertical smearing is present. High velocity anomalies are relatively better resolved and less affected by the smearing effects. Due to the smearing, the anomalous low velocity zones (<1.5 km/s) may be inversion artifacts due to the presence of a deeper low velocity layer which has smeared up to the surface. This effect is even stronger at greater depths where the velocity contrast is presumably higher. The model is not capable of resolving the abrupt boundary already discussed.

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Fig. 10. Synthetic resolution tests. Distance in km is calculated from the negative end of the Y-axis in Figs. 6 and 7, that is, an offset of 27 km is considered. Red circles indicate the instrument locations. (A) Synthetic velocity structure from the 3D model along the vertical slice along Y axis at X=0 km, which is roughly coincident with the central line L10. (B) Recovery results from the vertical slice shown in (A) masked according to the final velocity model. Note the significant smearing effects on the anomaly recovery. (C) Synthetic velocity structure from the 2D model along the central line L10 with an analogous geometry as in (A). (D) Recovery results from the 2D model shown in (C). Overall, the anomalies are better recovered in shape and amplitude than the 3D model results, albeit the transition zone shows some smearing.

To evaluate the robustness of our 3D model and improve some of its limitations along the central line, we also generated a 2-D velocity-depth model using the joint refraction/reflection 2-D tomographic inversion code of Korenaga et al. (2000) and additional a-priori information to generate a more complex starting model. For this analysis, we used the first arrival picks from the central profile L10 (Fig. 4), comprising a total of 2,920 picks (1,800 P_s phases and 1,120 P_g phases). Fig. 11A shows the 2D model results and Fig. 11B shows a velocity difference model between this 2-D model and the nearly coincident vertical model slice extracted from the 3-D model (Fig. 6C).

307 In general, the 2-D velocity structure along the central profile agrees well with the corresponding model 308 slice from the 3-D model results in the upper 2 km (Fig. 11B). This is not surprising since the 3D model shows 309 little variation across the X-direction that could bias the 2D interpretation of the 3D structure (Figs. 6, 7 and 8). 310 The primary model features discussed in the previous section are observed in the 2D model: consistently lower 311 shallow velocities in the southwest region and an increased vertical velocity gradient extending to the northeast 312 (Fig. 11A). Greater discrepancy in the upper 2 km is observed at ~30 km where the high velocity feature observed in both models is resolved with higher amplitude (>0.5 km/s) in the 2D model (Fig. 11B). The large discrepancy 313 314 at depth ($|DV_p| > 0.5$ km/s) reflects the smooth initial model and smoothing parameters used for the 3D model, 315 which do not resolve the abrupt velocity contrast suggested by the triplication observed in the data. For the 2D 316 modelling, the starting model was more realistic and the modelling approach allowed for a discontinuity. Fig. 11D 317 shows a good data fit of both phases and the position of the critical point.



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319 Fig. 11. Direct comparison of velocity-depth models along central line L10 (Figs. 6 and 7) of the study area and accuracy and 320 data fitting of the 2D model. Distance in km is calculated from the negative end of the Y-axis in Figs. 6 and 7, that is, an offset 321 of 27 km is considered with respect to Y-axis. (A) 2D velocity model obtained by travel time inversion of first-arriving P-322 waves collected along profile L10 using the joint refraction and reflection travel time tomography code of Korenaga et al., 323 (2000). (B) Velocity difference model between the 2D model and the final 3D velocity model along Y axis at X=0 km, roughly 324 coincident with central seismic profile L10. (C) 2D velocity uncertainty model for the final model shown in (B), using the 325 Monte Carlo analysis. Higher deviation values, up to ± 0.2 km/s (darker grey zones), are found at the bottom of the model and 326 as localized patches at the northeast. (D) Travel time fitting for the 2D model. Observed times and associated picking 327 uncertainty used for the 2D modelling (0.06 s for P_s phases and 0.08 s for P_g phases) are indicated by red and magenta dots 328 with their corresponding error bars, respectively. Calculated times (black dots) are obtained from the 2D final model shown 329 in (B). Final RMS data misfit are 53 and 69 ms for P_s and P_g phases, respectively. Poor fitting of P_g phase at OBS 1010 (at 330 $X \sim 37$ km) is consistent with higher uncertainties of the 2D model along the corresponding travel paths.

332	In order to assess the accuracy of the 2D seismic velocity model, we conducted a Monte Carlo uncertainty
333	analysis (Korenaga et al., 2000) in a similar fashion to the 3D case, to estimate the uncertainty of our model
334	parameters: our starting velocity model and smoothing constraints. This uncertainty parameter of the calculated
335	V_p is better constrained in the southwest region with most of the values lower than 0.01 km/s whereas most of the
336	northeast region values are higher than 0.02 km/s and several patches of increased values (up to \pm 0.1 km/s) are
337	present (Fig. 11C). Resolution tests show lower vertical smearing effects than in the 3D case (Fig. 10D). Poorer
338	recovery coincides with the position of the abrupt velocity gradient (Fig. 10D); however, it is very localized and
339	is correlated with comparatively higher uncertainty values (>0.02 km/s) (Fig. 11C). It also noteworthy that V_p <
340	1.5 km/s are not present in the 2D results (Figs. 11A, 11C).



342 Fig. 12. Tectonic interpretation and 3D map view sections shown in Fig. 8. (A) Experiment setting whose coordinate system 343 coincide with that of Figs. 6-8. Bathymetry and structural observations inferred by seismic reflection lines (coincident with 344 seismic refraction profiles from this study (Flueh and Bialas, 2008) and ENAP-1 (Geersen et al., 2011)). Fluid activity 345 indicators are also shown. The location of the thrust ridge and these fluid activity features are also overlain in map view 346 section at Z = 0 km, (B), coincident with the seafloor, for better understanding. Note that these accretionary features and fluid 347 activity related processes correlate with lower seismic velocities (< 2 km/s) while less perturbed sediment layers mainly 348 located to the northeast present relative higher seismic velocities (around 2 km/s). (C) Map view section at Z = 1 km depth, 349 although with less resolution indicates that the perturbed sedimentary area coincident with the thrust ridge shows lower 350 seismic velocities. (D) Map view at Z = 2 km agrees well with the aforementioned correlation and also indicate probable local 351 uplift along splay faults based on strong velocity contrasts.

352 **4. Discussion**

353 The three-dimensional results characterize the shallow sedimentary structure of the accretionary complex 354 in a transition zone to the continental framework rock based on the experiment location with respect to the 355 continental backstop (Fig.1-3). Comparison with 2D results along the central line shows that the velocity structure 356 is primarily two-dimensional (Fig. 11A, 11B) showing a spatial correlation between lower velocity zones in the 357 near-surface and the location of thrust ridges. Even though only shallow ray coverage is achieved, possible local 358 uplifts at ~30 km and ~40 km along profile L10 (Fig. 11A), are inferred by noteworthy lateral velocity gradients 359 (Fig. 12D). This feature correlates with the splay fault system that produces the observed thrust ridges and that 360 probably extends deeper within the framework rock (Geersen et al., 2011). Overall, the near-surface velocity 361 structure is better represented by the 2D model, which results in lower uncertainty values and improved resolution 362 due to finer grid spacing, a more accurate initial velocity model and less vertical smoothing. The 2D model also 363 shows less vertical smearing than the 3D model, which may explain the very low sedimentary velocities obtained 364 at the seafloor in the 3D model.

At shallow depths, the seafloor expression of the thrust ridges defines an area of comparatively lower velocities (< 1.7 km/s) (Fig. 11A, 12B) that extends to ~2 km beneath the sea surface (1 km below the seafloor) but with slightly increasing velocities due to sediment compaction (Fig. 12C). Many of these features are distributed over the ridges and the proximal intra-slope hemipelagic sedimentary basins that have been reported to be of biogenic origin, strongly diluted with high terrigenous input, mostly silty clays and clays (Flueh and Bialas, 2008) (Fig. 12C). There is also a spatial correlation between these zones and fluid activity inferred by heat flow anomalies, seafloor observations and sidescan mapping (Fig. 12).

Fig. 13A shows a seismic reflection line coincident with profile L10 that properly images the upper 500 m. A clear and continuous seafloor parallel reflection identified to be a gas hydrate-related bottom-simulating reflector (BSR), is mostly visible. However, several disruptions are observed and often coincide with dipping reflections, which might indicate a connection with the splay fault system mapped by Geersen et al., (2011) (Fig. 13B). The BSR indicates the base of the Gas Hydrate Stability Zone (GHSZ), which has a mean thickness of ~300 m along the profile. Because the BSR is caused by the impedance contrast between hydrate-bearing sediments and sediments containing free gas located right beneath it, it is a robust proxy for the presence of gas hydrate within
the GHSZ (Tréhu et al., 2003). Although the corresponding free gas zone beneath the GHSZ should form a low
velocity layer (Tréhu et al., 2001), this feature is not resolved by our models.

381 The near-surface velocity field within the GHSZ is not uniform along the profile (Fig. 13B). Even though 382 V_p values lower than 1.7 km/s are predominant in the vicinity of the thrust ridges, there are two zones where the 383 2.0 km/s velocity contour is found above or close to the BSR (at ~30 km and at the north-eastern end of the profile; 384 Fig. 13B). We speculate that decreasing V_p values within the GHSZ might indicate decreasing hydrate 385 concentration. Tréhu et al. (2001; 2004a) suggested that free gas presence within the GHSZ may contribute to the 386 local decrease in seismic velocity in the shallow sediment structure above the BSR. However, Tréhu et al. (2004b), 387 discussed the potential mechanisms that allow free gas stability within the GHSZ and proposed that high gas 388 saturation values may be the driving force for the focused flow to the GHSZ, that is, in focused conduits where 389 free gas can be isolated from the pore water, hindering gas hydrate formation. Although there are no observable 390 seismic indicators, such as bright spots above the BSR implying free gas conduits, (Fig. 13A), this situation cannot 391 be ruled out considering evidence of effective gas migration through the GHSZ, as will be discussed later.

392 Reported advective fluid flow inferred by high heat flow anomalies is observed in the thrust ridge vicinity 393 along profiles c728, ENAP-1 and VG06-74 in Figs. 12A and 12B (Grevemeyer et al., 2003; 2006; Villar-Muñoz 394 et al., 2013). These results have been postulated as a geothermal evidence of the migration of warm fluids into the 395 GHSZ through thrust faults that facilitate the vertical migration of fluids that originated at greater depth 396 (Grevemeyer et al., 2006; Villar-Muñoz et al., 2013). Fig. 13D shows heat flow values derived from the BSR 397 depths along profile L10, shown in Fig.13C (details of the calculations in the Supplementary material, section A3) 398 and also includes the corresponding values of the previously discussed independent results at their intersection 399 with this profile. The overall character of the results shows clear anomalies around fault-related reflectors for both 400 the lithostatic and hydrostatic cases. Thus, we will only focus on the heat flows derived for the hydrostatic case, 401 following the recommendation of Hyndman et al., (1992). Noteworthy shallowing of the BSR observed at 23 km 402 and 38 km (Fig. 13C) results in high heat flow values (~40 mWm⁻²; Fig. 11C) and are consistent with focused flow 403 along thrust faults (Fig. 11B; Tréhu et al., 2003). The profile section between 28 and 39 km presents comparatively 404 high heat flow rates (>35 mWm⁻²) roughly consistent with the independent BSR-derived values (Fig. 11D) and 405 above the theoretical thermal model at this distance from the deformation front reported by Grevemeyer et al., 406 (2003) (~30 mWm⁻² at ~40 km). Within this section, a gap between 30 and 32 km where no clear BSR is found 407 coincides with a very high direct heat flow measurement (~120 mWm⁻²; Grevemeyer et al., 2006) and a local uplift

- 408 inferred by the calculated velocity structure and associated with lithological discontinuities (see velocity contours
- 409 and inferred thrust faults in Fig. 13B).





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Lithostatic (Vp=1.8 km/s)

Villar-Muñoz et al., 2014 (BSR-derived/ Hydrostatic)

411 Fig. 13. (A) Depth-migrated seismic reflection coincident with profile L10. Conversion to depth was made using the 2D 412 model shown in 11A. (B) Depth-migrated seismic reflection profile shown in (A), overlain by 2D velocity grid from Fig. 413 11A. Some velocity contours from the 2D model and instrument locations are overlain. The presence of BSR is indicated by 414 red dots (also shown in (C)). Oblique reflections offsetting the BSR are interpreted as thrust faults (thick yellow lines). Thrust 415 ridge location is visible as local seafloor uplift and distinctive reflection signature characterized by deformed sedimentary 416 strata. In contrast, hemipelagic sediments in intra-slope basins are distinguished by predominantly parallel reflectors along a 417 smoother seafloor. (C) Selection of seafloor and BSR reflections digitized from the seismic image in (A). Faint reflections 418 and local disappearance of the BSR produce several gaps. Note disturbed BSRs around faults. (D) Heat flow values derived 419 from BSR depths shown in (C) for a hydrostatic case and lithostatic case with different mean velocity values for the hydrate-420 bearing sediments. Different independent measurements coincident with profile L10 are overlain.

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422 The presence of cold seeps has been documented by Klaucke et al., (2012) (Figs. 3 and 12). Seeps transport 423 dissolved and gaseous compounds upward and sustain oasis-type ecosystems at the seafloor (Suess et al., 2018) 424 such as the observed seep fauna comprising bacterial mats, shells of vesicolamyd clams (Archivesica sp) and large 425 tubeworms of the genus Lamellibrachia (Klaucke et al., 2012) (Figs. 3 and 12). Moreover, methane related 426 reactions -mainly AOM, anoxic oxidation of methane- involve consequent carbonate mineral precipitation near 427 the seafloor (Suess et al., 2018), e.g., in the form of carbonate buildups (chemoherms), observed by Klaucke et al., 428 (2012) (Figs. 3 and 12). As Suess et al. (2014) pointed out, cold seeps and their products have also been explained 429 as the result of tectonic fluid expulsion by dewatering of sediments. Dewatering occurs in response to lateral 430 compression by plate convergence and, as Cloos (1984) proposed, seafloor fluid flow through sediment facies may 431 be explained by thrust faults acting as dewatering conduits. Landward of the accretionary wedge, splay faults 432 developed in shallow angle subduction zones as off Cascadia, Japan and Southern Chile (as is the case of this 433 study) are also proposed as drainage pathways of the upper plate (Moore et al., 2007).

Klaucke et al. (2012) reported little evidence of current seepage activity as e.g., gas flares or bubbles in the water column, suggesting either intense seepage in the past or moderate activity over prolonged periods of time. According to velocity-derived porosities for the upper 2.5 km (more details on the calculations are included in the Supplementary material, section A4), the low velocity sediments (< 1.7 km/s) that spatially correlate with these seep sites show high porosity values (> 50%). Therefore, as Bangs et al. (1990) proposed for the case of Barbados Accretionary complex, an overpressure caused by a considerable amount of pore fluids coming from underlying sediment facies as a result of a dewatering process, may also partially account for a decreasing seismic velocity in the shallow sediments.



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443 Fig. 14. Schematic interpretation of the study area based on the seismic structure along profile L10 in Fig. 13, summarizing 444 the model results shown in Figs. 6-8, 11 and the tectonic interpretation in Fig. 12. Clusters of lower seismic velocity within 445 the GHSZ (<1.7 km/s, see Fig. 13B) might be interpreted as hydrate-bearing sediments containing less hydrate concentration. 446 Hemipelagic sediments in the intra-slope basins around the thrust ridges constitute the slope apron and are prone to experience 447 seep activity (seepages, carbonate formations, seep biota). The interpreted accretionary prism is comprised between the 2.0 448 and 4.0 km/s; however, a transition zone composed of moderately porous accreted sediments (40-50% porosity) is present as 449 noted by the region between the base of the GHSZ and the 2.0 km/s velocity contour (dotted grey line, see Fig. 13B for the 450 velocity contour reference). This boundary also implies localized tectonic uplift associated with the splay faults. Fluid 451 migration along these faults, represented by vertical blue arrows, may indicate the dewatering of the accretionary complex 452 implying a tectonically induced cold seep activity. Subsequent lateral fluid migration might also explain the existence of other453 adjacent reduced velocity zones within the GHSZ.

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Accordingly, we propose that the observed seeps, evidenced by carbonate formations and seep biota at the seafloor are explained by the upward migration of methane-rich warm fluids into the high porosity sediments within the GHSZ, through splay faults associated with the thrust ridges (Fig. 14). Fluid flow through the gas hydrate stability field may occur where warm fluids move relatively quickly through permeable settings and keep the surrounding sediments warm enough to prevent gas hydrate formation (Grevemeyer et al., 2006). A differential hydrate concentration might be revealed by the variation of the near-surface velocity structure within the GHSZ, suggesting probable lateral fluid migration within the GHSZ along inclined sedimentary strata (Fig. 14).

462 On the other hand, at the NW region of the model, comparatively higher shallow velocities (>1.7 km/s), lower porosities (<50%) and higher velocity-depth gradients ($\sim1.0 \text{ s}^{-1}$) are observed in areas of less disturbed 463 464 seafloor, corresponding to parallel and horizontally stratified hemipelagic sediment sequences (Flueh and Bialas, 465 2008). This shallow behaviour is observed in seismic refraction profiles at these distances from the deformation 466 front, outside the area of thrust ridges, as seen along P09 profile (see Fig. 2 for location). Even though resolution 467 at the edge of the model decreases, extrapolation to seismic velocities observed in P09 (Contreras-Reves et al., 468 2016) suggests that the increase in velocity is due mainly to the rock compaction processes unaffected by seepage 469 activity.

486 **5.** Conclusions

487 A shallow 3-D P-wave tomography of the landward edge of the accretionary prism offshore SC Chile 488 (35°- 36°S) has been determined. Results show a primarily 2-D velocity distribution characterized by the presence 489 of very low velocity zones within the GHSZ (<1.7 km/s), which are interpreted as highly porous sediments (higher 490 than 50%) hosting less gas hydrate. A spatial correlation between these zones and evidence of seep activity (seep 491 biota and BSR-derived heat flow anomalies as high as 40 Mwm⁻²) in the vicinity of the locally NS trending thrust 492 ridge indicate an increased supply of methane-rich fluids in the shallow sediments into the GHSZ. This situation 493 responds to the existence of an accretionary complex experiencing a dewatering process due to overburden and 494 tectonic stress.

Accretionary thrust ridges in this transition zone between the accretionary complex and the framework rock are the shallow expression of splay faults and are present in other accretionary margins such as the Nankai accretionary margin in Japan. In our study, these features are clear by their seafloor morphology signature and associated tectonic uplift has been inferred by increased lateral velocity gradients at depth. Thus, the dewatering of the accretionary complex is tectonically driven by these existing splay faults that act as fluid pathways for the draining of the upper plate.

The seismic velocity imaging of the upper sedimentary structure around thrust ridges can be a useful tool for the assessment of cold seep activity. However, better resolution needs to be achieved in order to characterize the GHSZ with more detail and additional seepage proxies are necessary for further insights into the relationship between shallow fluid activity in underconsolidated sediment deposits and their elastic properties such as V_p velocities.

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514 **7. Data availability**

515 The research data used in this work is publicly available in an OSF repository (Obando-Orrego et al.,

516 2020). It consists of SEG-Y files of each of the nine used instruments (hydrophone component) for the 5

- 517 seismic profiles (L10, L11, L12, L13, L14) and the corresponding navigation files (*ukooa* files). Also included
- are the corresponding migrated seismic reflection lines.

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676 Supplementary material

677 A1. Aftershock seismicity





684 A2. 3D seismic records

685 The following figure shows examples of the seismic records from the instrument OBH 1002 along distal686 lines, that is, L11, L12, L13 and L14 lines (see Fig. 3 for profile locations).





688

Supplementary Figure 2. Seismic records from instrument OBH 1002 along (A) L11, (B) L12, (C) L13, (D) L14 (see Fig. 3 in the main text for map reference). For each seismic record, the detected first-arriving P waves refractions are indicated: P_w (water wave arrival), P_s (low velocity sedimentary phase) and P_g (deeper sedimentary phases). Traces are ordered by shot number.

Heat flow values along profile L10 derived by BSR depths, using a simple linear conductive heatrelationship:

696
$$q = k (T_{BSR} - T_{seafloor}) / (Z_{BSR} - Z_{seafloor})$$

697 where q is the heat flow, k is the thermal conductivity, $T_{seafloor}$ is the temperature at the seafloor, T_{BSR} is the 698 temperature at the BSR, Z_{BSR} and $Z_{seafloor}$ are depths of the BSR and seafloor, respectively. Thermal conductivity 699 and seafloor temperature were obtained from Grevemeyer et al., (2003) based on ODP Leg 202 drillcore data 697 (k=0.85 W/mK; $T_{seafloor}=5^{\circ}$ C). The BSR and the seafloor depths were obtained from the seismic profile shown in 708 Fig. 11A. Z according to seismic profile. Temperature at BSR depth T_{BSR} is calculated using the dissociation 709 temperature-pressure function T(p) of gas hydrates published by Dickens and Quinby-Hunt, (1994) for a seawater-709 methane system:

704
$$1/T = 3.79 \cdot 10^{-3} - 2.83 \cdot 10^{-4} \log(P)$$

where p is the pressure at the BSR (MPa) and T the temperature (Kelvin). Gas in the system is assumed to be pure methane, with a pore water salinity of 35 g 1^{-1} . The pressure at BSR depth is studied for two cases: hydrostatic and lithostatic equilibrium. A density value for seawater of 1020 kg/m³ was used and density for the sediment column were calculated using the relationship for soft sediments from Hamilton et al., (1978):

709
$$\rho = 1.135 \text{ V} - 0.19$$

where V is velocity in km/s. We explored solutions for V = 1.6, 1.7 and 1.8 km/s for the hydrate-bearing sediments.

712	In order to assess the degree of pore pressure in the shallow sediments, we calculated velocity-derived
713	porosity values of the upper 2 km (Vp < 2.5 km/s). Empirical density-porosity and velocity-porosity relationships
714	for characteristic deep-sea sediments have been derived by Hamilton (1978) and Hyndman et al., (1993),
715	respectively. Hyndman's relation was derived from laboratory experiments of marine sediments in the Nankai
716	accretionary margin, Japan, and by first approximation, was applied here in the valid range of porosities of 30-
717	60% (~1600 and 2500 m/s). According to this relation,
718	
719	(1) $P = -1.180 + 8.607 (1/V) - 17.89 (1/V)^2 + 13.94 (1/V)^3$
720	
721	where P is porosity (%) and V is velocity in km/s.
722	
723	Alternatively, we computed densities from seismic velocities using Hamilton's relation derived from
724	carbonate silt clays, turbidites, mudstones, and shales forming the soft and unlithified upper 500 m sediment layers,
725	which is valid in the approximate range of 1.53-2.0 km/s:
726	
727	(2) $\rho = 1.135 \text{ V} - 0.19,$
728	
729	where V is the P-wave velocity (km/s) and ρ is density in g/cm ³ . Porosity values were then derived using the
730	corresponding definition:
731	
732	(3) $P = 100 ((\rho - \rho_m) / (\rho_f - \rho_m)),$
733	
734	where P is porosity (%), ρ is density (g/cm3), ρ_f corresponds to seawater density of 1.020 g/cm3 and ρ_m is the grain
735	density. We used a value of 2.7 g/cm ³ corresponding to silica.



737 Supplementary Figure 3. Porosity versus velocity using Hamilton and Hyndman relations. Dashed sections of the curves
738 indicate extrapolation of the porosity values outside the respective valid ranges. Accordingly, velocities lower than 1.53 km/s
739 might imply porosities higher than 70%.









Legend



Seep fauna (Klaucke et al., 2012)





















-4

-2

2

X [km]















(8)

(D)

.



(C)

Declaration of interests

 \boxtimes The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

□ The authors declare the following financial interests/personal relationships which may be considered as potential competing interests:

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